Numerical Simulations of a Delay Model for Immune System-Tumor Interaction

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ABSTRACT: In this paper we consider a system of delay differential equations as a model for the dynamics of tumorimmune system interaction. We carry out a stability analysis of the proposed model. In particular, we show that the system can have up to two steady states: the tumor free steady state, which always exist, and the tumor persistent steady state, which exists only when the relative rate of increase of the tumor cells exceeds the ratio between the natural proliferation rate and the relative death rate of the effector cells. We also determine an upper bound for the delay, such that stability is preserved. Numerical simulations of the system under different parameter values are performed.

Keywords: Delay differential equations, Asymptotic stability, Numerical simulations.

المحاكاة العددية لنموذج تفاعل نظام المناعة والورم السرطاني باستخدام المعادلات التفاضلية التأخيرية

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الملخص: نفترض في هذا البحث جملة من المعادلات التفاضلية التأخيرية كنموذج لديناميكيات نظام التفاعل بين الأورام السرطانية والجهاز المناعي. ونقوم بتحليل استقرار النموذج المقترح. كما نبين بشكل خاص امكانية احتواء النظام على حالتين ثابتتين: الحالة المستقرة الخالية من الورم، وهي موجودة دائمًا، والحالة المستقرة للورم المستمر، والموجودة فقط عند تجاوز المعدل النسبي لزيادة الخلايا الورمية النسبية بين معدل الانتشار الطبيعي ومعدل الوفيات النسبي للخلايا المستجيبه. ونعين أيضا الحد الأعلى للتأخير مع المحافظة على الاستقرار. وتم إجراء محاكمة عديماً ال

الكلمات المفتاحية: معادلات التفاضلية التأخيرية، تقارب الاستقرار، محاكاة عددية.

1. Introduction

The immune system (IS) is a very complex one. It is composed of a complex network of different cells which operate collectively by communicating information through signalling. The duty of these cells is to keep our system natural, i.e., free of foreign entities. When a foreign entity enters our system, particular immune system cells raise an alarm and send signals through the network to other immune system cells calling for an attack on the foreign entity. Once the foreign substance is located and identified as non-self, the immune system plans an attack strategy. Among the most deadly foreign entities are cancer cells. Cancer cells are difficult to deal with because of their ability to multiply so fast that the IS cannot keep up. Another property of cancer cells is that they are able to camouflage the

antigens (the substance that triggers the IS) such that the IS cannot recognise them as foreign cells.
Once cancer cells are identified as foreign, the IS starts an attack to destroy them. The interactions between tumour cells and the immune system are very complex and require sophisticated models to describe them.
Mathematical models, based on ordinary differential equations, delay differential equations or partial differential equations. Several mathematical models have been suggested to describe the interactions between tumour and immune system [1-9].



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Many of these papers consider the tumor-IS interaction with other factors such as treatment. In this paper, we consider a mathematical model of IS-tumor interactions based on a system of delay differential equations. The delay is introduced to reflect the non-instantaneous outcome of the interaction between the immune system cells and the tumor cells. We carry out a stability analysis of the system and we simulate the system under different parameter values to explore the different asymptotic behaviors.

Many models that have appeared in the literature are based of the well-known Kuznetsov and Taylor model [1] which we describe in the next section.

2. Kuznetsov and Taylor model

The Kuznetsov and Taylor mathematical model to describe the tumor-immune system interaction is based on the following:

1. The growth of tumor cells population T (in the absence of the immune system cells) follows a logistic model:

$$\frac{dT}{dt} = aT \left(1 - bT_{tot}\right),\tag{1}$$

where *a* is the maximum growth rate, b^{-1} is the carrying capacity of the biological environment for the tumor cells, and T_{tot} is the total population of unhit *T* cells, $T_{tot} = T + C$.

2. The rate of production of the immune system cells E (the effector cells), has two sources: (i) a constant normal production rate s (in the absence of the tumor cells) and (ii) a production rate caused by the presence of tumor cells. The magnitude of the second production rate is a function, F(C,T), of the current concentrations of C and T, where C is the concentration of the complex E-T. The immune system cells E die out at a rate d_1 .

3. The interaction between the effector cells E and the tumor cells T forming the complexes $C \equiv E - T$ is described by the following kinetic

$$E+T \xleftarrow{k_1}{k_{-1}} C \xleftarrow{k_2}{k_3} E+D_T$$
$$D_E+T$$

Figure 1. Interaction kinetics between E and T cells.

where the C complex is formed at rate of k_1 and breaks at three different rates: k_{-1} giving E and T unaltered, k_2 given E (unaltered) and D_T , and k_3 giving D_E and T (unaltered), where D_T and D_E are altered T and E cells which are bound to die.

Kuznetsov and Taylor's complete model is given by the following set of differential equations

$$\frac{dE}{dt} = s + F(C,T) - d_1 E - k_1 ET + (k_{-1} + k_2)C,$$

$$\frac{dT}{dt} = aT (1 - bT_{tot}) - k_1 ET + (k_{-1} + k_3)C,$$

$$\frac{dC}{dt} = k_1 ET - (k_{-1} + k_2 + k_3)C,$$

$$\frac{dD_E}{dt} = k_3 C - d_2 D_E,$$

$$\frac{dD_T}{dt} = k_2 C - d_3 D_T,$$
(2)

where, again, E, T, C, D_E , D_T are the concentrations at the tumor site of the immune system cells (effector cells), the tumor cells, the effector-tumor cells complex, the inactived effector cells, and the inactived tumor cells, respectively. The above model was simplified based on the fact that the E-T bond, which is formed at a rate of k_1 , lasts a relatively short period of time [10] and disassociates giving rise to E and T (no win situation) with a rate k_{-1} ,

or D_E and T (tumor wins) with a rate k_3 , or D_T and E (E wins) with a rate k_2 , see Figure 1. The inactivated effector cells D_E and inactivated tumor cells D_T are bound to die with rates d_2 and d_3 , respectively.

In [1], it was claimed that C is approximately constant, that is $\frac{dC}{dt} \approx 0$, hence $C \approx KET$, with $K = k_1/(k_{-1} + k_2 + k_3)$. In addition, T_{tot} was approximated by $T_{tot} \approx T$, and $F(C,T) \equiv F(E,T)$.

These assumptions reduce the above system (2) to the following IS-tumor interaction model, in non-dimensional form,

$$\frac{dE}{dt} = s + F(E,T) - mET - dE,$$

$$\frac{dT}{dt} = aT(1 - bT) - nET,$$
(3)

where the parameters $m = Kk_3$, $n = Kk_3$, and $d = d_1$. The function F(E,T) describes the rate at which cytotoxic effector cells accumulate around the region where the T cells are localized. Based on Kuznetsov and Taylor's model, many researchers have derived other models by considering different forms for the function F(E,T) and/or by adding additional considerations (see [3]-[8], and references therein).

In [1], the authors analyzed model (3) with a Holling Type II form for F. In non-dimensionly variables, their model takes the form

$$\frac{dx}{dt} = \sigma + \frac{\rho xy}{\eta + y} - \mu xy - \delta x,$$

$$\frac{dy}{dt} = \alpha y (1 - \beta y) - xy.$$
 (4)

Later, in [6], the authors considered Kuznetsov and Taylor's model (3) with Holling Type I response function F(E,T) = kET, with and without delay, where the delay was introduced to reflect the delay in the response of the immune system before proliferating effector cells.

In this work, we consider (3) with a Holling Type I response function F(E,T) = kET with delay introduced in both equations as described in the next section.

3. A delay model

As in [6], we consider F(E,T) of Holling Type I, but we consider the delay τ according to the following biologically possible scenario. We assume that there is delay in both the proliferation of effector cells and in the disposal of tumor cells. In dimensionless form, the model we consider is

$$x'(t) = \sigma + wx (t - \tau) y (t - \tau) - \delta x,$$

$$y'(t) = \alpha y (1 - \beta y) - x (t - \tau) y (t - \tau),$$
(5)

where x(t) and y(t) are the nondimensional density at the tumor site of the effector cells and of the tumor cells, respectively. All parameters in (5) are nonnegative real numbers and are approximated (see [1] and [6]) by:

$$\sigma = 0.1181, \qquad w = 0.04, \qquad \delta = 0.3743, \alpha = 1.636, \qquad \beta = 2.0 \times 10^{-3}.$$
(6)

The delay factor in model (5) models the natural delays in both the proliferation of effector cells and in the disposal of tumor cells. This is biologically meaningful as the IS takes time to react and migrate effector cells to the tumor site. Also, in the second equation of model (5), it models the delay in completely eradicating the tumor cells as they will not die immediately upon interacting with effector cells.

Our aim in studying the above model (5) is to investigate the effect, if any, of the time delay, τ , in the disposal of the tumor cells. We carry out the stability analysis of the equilibrium states of the model (5) and determine necessary conditions for the local asymptotic stability. Also, we determine an upper bound for the delay parameter τ such that stability is preserved. Numerical simulations of the model are performed to investigate the asymptotic behavior of the system.

4. Steady states

The steady states, (\bar{x}, \bar{y}) , of (5) are the solutions of the nonlinear homogeneous system

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$$\sigma + w\overline{xy} - \delta \overline{x} = 0,$$

$$\alpha \overline{y} (1 - \beta \overline{y}) - \overline{xy} = 0.$$
(7)

From the second equation of (7) we have either $\overline{y} = 0$ or $\overline{x} = \alpha(1 - \beta \overline{y})$. If $\overline{y} = 0$, then from the first equation of (7) we have $\overline{x} = \sigma/\delta$, giving the tumor free steady state $E_0 = (\sigma/\delta, 0)$. If $\overline{y} \neq 0$, then $\overline{x} = \alpha(1 - \beta \overline{y})$ and from the first equation of (7) we have the quadratic equation for \overline{y} :

$$\alpha\beta w \overline{y}^2 - \alpha(\beta\delta + w)\overline{y} + \alpha\delta - \sigma = 0$$

whose discriminant Δ is

$$\Delta = \alpha^2 (\beta \delta - w)^2 + 4\alpha \beta \sigma w > 0.$$

Then depending on the parameter values, there may exist up to two more steady states $E_1(x_1, y_1)$ and $E_2(x_2, y_2)$, where

$$x_{1} = \frac{-\alpha(\beta\delta - w) + \sqrt{\Delta}}{2w}, \qquad y_{1} = \frac{\alpha(\beta\delta + w) - \sqrt{\Delta}}{2\alpha\beta w}$$
(8)

$$x_{2} = \frac{-\alpha(\beta\delta - w) - \sqrt{\Delta}}{2w}, \qquad y_{2} = \frac{\alpha(\beta\delta + w) + \sqrt{\Delta}}{2\alpha\beta w}$$
(9)

The existence of E_1 and E_2 as biologically meaningful (i.e., with positive coordinates) depends on the parameters. If, as assumed, all parameters are positive, then it is clear that $x_2 < 0$ because $\sqrt{\Delta} > |\alpha(\beta\delta - w)|$. So E_2 does not exist as biologically meaningful. As for E_1 , since $\sqrt{\Delta} > |\alpha(\beta\delta - w)|$, we have $x_1 > 0$. It remains to check when $y_1 > 0$. Calculations reveal that $y_1 > 0$ if and only if $\sigma/\delta < \alpha$. Therefore, if $\alpha < \sigma/\delta$, we have only the tumor free steady state $E_0 = (\sigma/\delta, 0)$ and if $\sigma/\delta < \alpha$, we have two steady states $E_0 = (\sigma/\delta, 0)$ and $E_1 = (x_1, y_1)$, where x_1 and y_1 are given by (8).

The value σ/δ relative to α is very important as it determines how many steady states the system can asymptotically assume. The value σ/δ measures the rate of increase of the density of the effector cells and α measures the rate of increase of the tumor cells. Thus if $(\sigma/\delta) > \alpha$, this biologically means that the effector cells will take over the tumor cells, which as a result, are bound to die in the long run. On the other hand, if $(\sigma/\delta) < \alpha$, the tumor cells will always exist.

Now, we turn to the stability analysis of the steady states E_0 and E_1 .

5. Stability analysis

We know that a steady state is locally asymptotically stable if all the eigenvalues of the Jacobian matrix of the system at the steady state have negative real parts.

Before we continue, we recall an important stability criterion known the Mikhailov criterion [9]:

Mikhailov criterion: Let P(s) and Q(s) be two polynomials with deg(P) = N and deg(Q) < N. If the quasi-polynomial $R(s) = P(s) + Q(s)e^{-s\tau}$ has no roots on the imaginary axis, then all roots of R(s) have negative real parts if and only if the argument of R(js), $j = \sqrt{-1}$, increases by $N \pi/2$ as $s \to +\infty$. For the system (5) under consideration, the Jacobian matrix at a steady state (\bar{x}, \bar{y}) is given by

$$J = \begin{bmatrix} -\delta + w\bar{y}e^{-\lambda\tau} & w\bar{x}e^{-\lambda\tau} \\ -\bar{y}e^{-\lambda\tau} & \alpha - 2\alpha\beta\bar{y} - \bar{x}e^{-\lambda\tau} \end{bmatrix}.$$
 (10)

The characteristic polynomial is

$$W(\lambda) = P(\lambda) + Q(\lambda)e^{-\lambda\tau}, \qquad (11)$$

where

$$P(\lambda) = \lambda^2 + (\delta + A)\lambda + \delta A, \qquad A = 2\alpha\beta\overline{y} - \alpha, \tag{11}$$

$$Q(\lambda) = (\delta \overline{x} - w \overline{y} A) + \lambda (\overline{x} - w \overline{y}).$$
⁽¹²⁾

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Stability of the tumor-free steady state E_0

For the steady state $E_0 = (\sigma/\delta, 0)$, the characteristic polynomial reduces to

$$W(\lambda) = (\lambda + \delta)(\lambda - \alpha + \sigma / \delta e^{-\lambda \tau}).$$
⁽¹³⁾

We have the following proposition.

Proposition 1

(i) If $\alpha > \sigma/\delta$, then E_0 is an unstable saddle point.

(ii) If $\alpha < \sigma / \delta$, then there exists a $\tau_0 > 0$ such that E_0 is locally asymptotically stable for $\tau \in [0, \tau_0)$ and unstable for $\tau > \tau_0$ where

$$\tau_0 = \frac{\cos^{-1}(\alpha\delta/\sigma)}{\sqrt{\sigma^2/\delta^2 - \alpha^2}}.$$
(14)

Proof. From (14), we see that one eigenvalue is $\lambda = -\delta < 0$. Next, we need to analyze the roots of the quasipolynomial

$$R(\lambda) = \lambda - \alpha + \sigma / \delta e^{-\lambda \tau}.$$
(15)

It is clear that if $\alpha > \sigma/\delta$, $R(\lambda)$ in (5.7) has a positive root, since the functions $f(\lambda) = \alpha - \lambda$ and $g(\lambda) = \sigma/\delta e^{-\lambda \tau}$ satisfy $f(0) = \alpha > g(0) = \sigma/\delta$ and $f(\lambda) \to -\infty$ and $g(\lambda) \to 0$, as $\lambda \to \infty$. This proves (i).

For (ii), assume that $\alpha < \sigma / \delta$. According the Mikhailov criterion, if $R(\lambda)$ has no imaginary roots, then all its roots have negative real parts if and only if

$$\left[Arg(R(js))\right]_{s=0}^{s=\infty} = \pi/2$$

Let $\lambda = js$, $s \in \mathbb{R}^+$, $j = \sqrt{-1}$. Then (16) becomes $R(is) = [-\alpha + \sigma/\delta co$

 $R(js) = [-\alpha + \sigma / \delta \cos(s\tau)] + j[s - \sigma / \delta \sin(s\tau)].$

Let $\theta = Arg(R(js))$. Then as $s \to \infty$, we have

$$\cos(\theta) \to 0$$
 and $\sin(\theta) \to 1$. (17)

(16)

This implies that $\lim_{s\to\infty} Arg(R(js)) = \pi/2$. We have $R(0) = -\alpha + \sigma/\delta > 0$ (for $\alpha < \sigma/\delta$) and Arg(R(0)) = 0. Thus

$$\left[Arg(R(js))\right]_{s=0}^{s=\infty} = \pi/2$$

which implies, according the Mikhailov criterion, that all roots of $R(\lambda)$ have negative real parts, provided that $R(\lambda)$ has no imaginary roots. Now let us see under what condition $R(\lambda)$ has imaginary roots. If R(js) = 0 for some $s \in R^+$, then

$$\alpha - \sigma / \delta \cos(s\tau) = 0, \tag{18}$$

$$s - \sigma / \delta \sin(s\tau) = 0. \tag{19}$$

Squaring and adding (19) and (20) gives

 $\underline{s^2} + \alpha^2 = \sigma^2 / \delta^2 \Longrightarrow s = \sqrt{\sigma^2 / \delta^2 - \alpha^2}.$

Conversely, if $s = s_0 = \sqrt{\sigma^2/\delta^2 - \alpha^2}$, the first value of τ , say τ_0 , that satisfies both (19) and (20) is

$$\tau_0 = \frac{1}{s_0} \cos^{-1}(\alpha \delta / \sigma) = \frac{\cos^{-1}(\alpha \delta / \sigma)}{\sqrt{\sigma^2 / \delta^2 - \alpha^2}}.$$
(20)

This implies that for $\tau < \tau_0$, $R(\lambda)$ has no imaginary roots, and E_0 is locally asymptotically stable. This completes the proof of Proposition 1.

Stability of the steady state $E_1 = (x_1, y_1)$

For $E_1 = (x_1, y_1)$, from equations (11)-(13), the characteristic polynomial takes the form

$$W(\lambda) = P(\lambda) + Q(\lambda)e^{-\lambda\tau}, \qquad (21)$$

where

$$P(\lambda) = \lambda^{2} + (\delta + A)\lambda + \delta A, \qquad Q(\lambda) = (C_{0} - \delta A) + \lambda(x_{1} - wy_{1}), \tag{22}$$

where

$$A = 2\alpha\beta y_1 - \alpha, \qquad C_0 = \delta A + \delta x_1 - w y_1 A.$$

Before we proceed, we state and prove the following Lemma about the positivity of C_0 , which will be used later.

Lemma 2 $C_0 = \delta A + \delta x_1 - wy_1 A = y_1 \sqrt{\Delta} > 0$. *Proof.* From the expression of $A = 2\alpha\beta y_1 - \alpha$ and $x_1 = \alpha(1 - \beta y_1)$, we can write A as $A = \alpha\beta y_1 - x_1$. Substituting this into C_0 , we get

$$C_0 = y_1 [\alpha \beta \delta - w(\alpha \beta y_1 - x_1)].$$

Now, from the expressions of x_1 and y_1 (Eq. (8)), we have

$$\alpha\beta y_1 - x_1 = \frac{\alpha(\beta\delta + w) - \sqrt{\Delta}}{2w} - \frac{-\alpha(\beta\delta - w) + \sqrt{\Delta}}{2w} = \frac{\alpha\beta\delta - \sqrt{\Delta}}{w}$$

It follows that

$$C_0 = y_1[\alpha\beta\delta - w(\alpha\beta y_1 - x_1)] = y_1\sqrt{\Delta} > 0.$$

We have the following proposition for the stability of E_1 .

Proposition 3. Assume that $C_0 > 2\delta A$. The steady state $E_1 = (x_1, y_1)$ is locally asymptotically stable, for $\tau \in [0, \tau_0)$, where τ_0 is given by (24).

Proof. Let $\lambda = js, s \in \mathbb{R}^+$. We have

$$W(js) = P(js) + Q(js)e^{-js\tau} = X(s) + jY(s)$$
(23)

where

$$X(s) = (-s^{2} + \delta A) + (\delta x_{1} - wy_{1}A)\cos(s\tau) + s(x_{1} - wy_{1})\sin(s\tau),$$
(24)

$$Y(s) = s(\delta + A) + s(x_1 - wy_1)\cos(s\tau) - (\delta x_1 - wy_1 A)\sin(s\tau).$$
 (25)

If $\theta = Arg(W(js))$, it can be easily verified that

$$\cos(\theta) = \frac{X(s)}{|W(j\omega)|_{w \to \infty}} - 1, \qquad \sin(\theta) = \frac{Y(s)}{|W(j\omega)|_{w \to \infty}} 0.$$
(26)

This implies that $\lim_{s\to\infty} Arg(W(js)) = \pi$. For s = 0, we have

$$W(0) = \delta x_1 - w y_1 A + \delta A = C_0 > 0,$$
⁽²⁷⁾

which implies that Arg(W(0)) = 0. Hence, $[Arg(W(js))]_{s=0}^{s=\infty} = \pi$, and E_1 would be locally asymptotically stable, provided that $W(\lambda)$ has no roots on the imaginary axis.

Now, we check for pure imaginary roots of $W(\lambda)$. If there exists $s \in R^+$ such that W(js) = 0, then X(s) = Y(s) = 0, or

$$a_{11}\cos(s\tau) + sa_{12}\sin(s\tau) = (s^2 - \delta A),$$
(28)

$$sa_{12}\cos(s\tau) - a_{11}\sin(s\tau) = -s(\delta + A),$$
 (29)

where

$$a_{11} = \delta x_1 - w y_1 A, \qquad a_{12} = x_1 - w y_1.$$

Squaring both sides of (29) and (30) and adding, we get

$$s^4 + a_1 s^2 + a_2 = 0, (30)$$

where

$$a_1 = \delta^2 + A^2 - a_{12}^2, \qquad a_2 = (\delta A)^2 - a_{11}^2 = -C_0(C_0 - 2\delta A)$$

The solutions of (31) are formally written as

$$s^{2} = \frac{-a_{1} \pm \sqrt{a_{1}^{2} - 4a_{2}}}{2}.$$
(31)

It is clear that a sufficient condition for the right hand side of (32) to be positive is $a_2 < 0$, i.e., $C_0 > 2\delta A$, since $C_0 > 0$. In this case, the positive solution of (32) is

$$s_0 = \sqrt{\frac{-a_1 + \sqrt{a_1^2 - 4a_2}}{2}} \tag{32}$$

Therefore, if $C_0 > 2\delta A$, then $W(\lambda)$ has only two pure imaginary roots $\pm js_0$, where s_0 is given by (33). With s_0 as in (33), the first positive value of τ such that X(s) = 0 and Y(s) = 0 is

$$\tau_0 = \frac{1}{s_0} \cos^{-1} \left(\frac{(a_{11} - a_{12}(\delta + A))s_0^2 - a_{11}\delta A}{a_{11}^2 + s_0^2 a_{12}^2} \right).$$
(33)

Therefore, if $C_0 > 2\delta A$, then for $\tau < \tau_0$, $W(\lambda)$ has no imaginary roots, and E_1 is locally asymptotically stable. \Box

6. Numerical simulations

In this section, we perform a number of numerical simulations to reveal the dynamics of system (5). We consider the two cases (i) $\alpha < \sigma/\delta$ where only E_0 exists, and (ii) $\alpha > \sigma/\delta$ where both E_0 and E_1 exist.

Case 1: For the case $\alpha < \sigma/\delta$, we simulate the system with the values of $\alpha = 0.1$ and $\alpha = 0.3$. The other parameters are fixed as in (6). In this case, the only equilibrium state is $E_0 = (\sigma/\delta, 0) = (0.315522, 0)$ which was proved to be asymptotically stable up to an upper bound for the delay τ (see Proposition 1 and Eq. (21)). The calculated upper bounds according to (21) are $\tau_0 = 4.17134$ and $\tau_0 = 3.22237$ for $\alpha = 0.1$ and 0.3, respectively. We remark that these upper bounds are valid only when the solutions x(t) and y(t) remain non-negative. In fact, simulations reveal that the solutions become negative for much smaller values of τ , $\tau \approx 0.05$. The outcomes of the simulations for this case are displayed in Figures 2, 4, and 5. In Figure 2, we display the phase portraits of system (5) for $\alpha = 0.1$ and $\alpha = 0.3$ and $\tau = 0$ and $\tau = 0.05$, each with 5 differential initial conditions. They all confirm the stability of E_0 , as the solutions x(t) and y(t) asymptotically tend to E_0 . In Figures 4 and 5, the solutions of system (5) are displayed with $\alpha = 0.1$ and $\alpha = 0.3$, respectively, each with $\tau = 0.05$. Both figures confirm the convergence of the solutions to the steady state E_0 .

It is important to mention that for $\tau > 0.05$, under certain initial conditions, the solutions become negative at some time t^* at which the system becomes biologically invalid.

Case 2: For the case $\alpha > \sigma/\delta$, we fix all parameters as in (6) and simulate the system with different delay values τ . In this case, we have two steady states, $E_0 = (\sigma / \delta, 0) = (0.315522, 0)$ which is unstable and $E_1 = (x_1, y_1) = (1.61138, 7.52522)$ which was proven to be locally asymptotically stable for up to $\tau < \tau_0$. From Eq. (34), τ_0 is calculated to be $\tau_0 = 0.0880838$. To confirm these results, we simulated the system for $\tau = 0,0.01,0.05,0.08,0.09$ and various initial conditions. The results are displayed in Figures 6 and 7, where Figure 6 displays the phase portraits for $\tau = 0, 0.01, 0.05, 0.08, 0.09$ and Figure 7 displays the solutions of the system for $\tau = 0.05, 0.08, 0.09$ and initial condition x(0) = 1, y(0) = 30. It is to be noted that for $\tau = 0, 0.01, 0.05, 0.08 < \tau_0$, the solutions converge to E_1 , confirming the local asymptotic stability of E_1 . However, in the case $\tau = 0.9 > \tau_0$, we notice the formation of the closed orbit around E_1 . This also can be seen in the rightmost plot in Figure 7, where the effector and tumor cells oscillate around E_1 . This confirms the existence of a bifurcation with respect to the parameter τ as it crosses $\tau_0 = 0.0880838$.

It should be also mentioned that for certain initial conditions the solutions of the system become negative in finite time.



Figure 2. Phase portraits of system (5) with different initial conditions (x (0), y (0)) = (0,1), (1,5), (2,10), (4,15)and (5,20): (a) $\alpha = 0.1$ and $\tau = 0$, (b) $\alpha = 0.1$ and $\tau = 0.05$. The equilibrium point is $E_0 = (0.315522, 0)$.



Figure 3. Phase portraits of system (5) with different initial conditions (x (0), y (0)) = (0, 1), (1, 5), (2, 10), (4, 15)and (5, 20): (a) $\alpha = 0.3, \tau = 0$, (b) $\alpha = 0.3$ and $\tau = 0.05$. The equilibrium point is $E_0 = (0.315522, 0)$.



Figure 4. Solutions of system (5) with $\alpha = 0.1$ and $\tau = 0.05$ and different initial conditions: (a) x(0) = 0, y(0) = 1; (b) x(0) = 1, y(0) = 5; (c) x(0) = 2, y(0) = 10; (d) x(0) = 4, y(0) = 15.



Figure 5. Solutions of system (5) with $\alpha = 0.3$ and $\tau = 0.05$ and different initial conditions: (a) x(0) = 0, y(0) = 1; (b) x(0) = 1, y(0) = 5; (c) x(0) = 2, y(0) = 10; (d) x(0) = 4, y(0) = 15.



(b) $\alpha = \sigma / \delta = 0.1636$ and $\tau = 0.05$; (c) $\alpha = \sigma / \delta = 0.1636$ and $\tau = 0.08$; (d) $\alpha = \sigma / \delta = 0.1636$ and $\tau = 0.09$.



Figure 7. Solutions of system (5) with initial conditions: x(0) = 1, y(0) = 30: (a) $\alpha = \sigma/\delta = 0.1636$ and $\tau = 0.05$; (b) $\alpha = \sigma/\delta = 0.1636$ and $\tau = 0.08$; (c) $\alpha = \sigma/\delta = 0.1636$ and $\tau = 0.09$.

7. Conclusion

A mathematical model for the interaction between the immune system cells and tumor cells was proposed and studied. The model was derived from Kuznetsov and Taylor's model [1] by introducing a delay term which reflects the slow response of the immune system cells to the presence of tumor cells. The importance in studying this delay model is in determining the effect of the delay term on the dynamics of the system. A stability analysis of the equilibrium states was carried out. Numerical simulations of the system were done for a number of different initial conditions. The simulation results confirm the theoretical stability results. Numerically, it was also observed that as the delay term τ crosses the upper bound τ_0 , the steady state E_1 loses its stability, resulting in the formation of closed orbits around E_1 . This is an important result in the sense that as the IS gets weaker (slower response, i.e., larger delay), the size of the tumor cells persists in an oscillating fashion (closed orbits).

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