# OPTIMIZATION OF ANTI-SYMMETRICAL OPEN CROSS-SECTIONS OF COLD-FORMED THIN-WALLED BEAMS 

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The paper deals with cold-formed thin-walled beams with the Z-, S- and Clothoid-section. A short survey of optimal designs of thin-walled beams with open cross-sections is given. Geometric properties of three crosssections are described. Strength, local and global buckling conditions for thin-walled beams are presented. The optimal design criterion with a dimensionless objective function as a quality measure is defined. Results of numerical calculations for optimal shapes of three cross-sections are presented in tables and figures.
Key words: thin-walled beam, open cross-section, global and local buckling, optimal design

## Notations

| $a, b, c, d$ | - dimensions of cross-sections |
| :--- | :--- |
| $r$ | - radius of the circular arc |
| $t$ | - thickness of the beam wall |
| $u$ | - dimensionless parameter of the clothoid |
| $A$ | - area of the cross-section |
| $H$ | - depth of the beam |
| $L$ | - length of the beam |
| $J_{S-V}$ | - geometric stiffness for Saint-Venant torsion |
| $J_{y}, J_{z}$ | - inertia moments |

$J_{\omega} \quad-\quad$ warping moment of inertia
$M_{0} \quad-\quad$ loading moment
$R \quad-\quad$ principal radius of the clothoid
$\alpha, \beta-$ angles of the S-section
$\lambda-$ relative length of the beam
$\theta_{p} \quad-\quad$ angle to the principal axes
$\omega \quad$ - warping function
$\Phi_{j} \quad-\quad$ objective function

## 1. Introduction

Shapes of open cross-sections of contemporary cold-formed thin-walled beams are rather complicated. They are usually mono-symmetrical, although sometimes anti-symmetrical too. The main constraints in designing thin-walled structures are strength and stability conditions. The beginnings of the optimal design of thin-walled structures reach back to 1959. The first paper on optimal design of a thin-walled beam with an open cross-section (I-section) in pure bending state was presented by Krishnan and Shetty (1959). A complete survey of optimal design problems of structures for the second half of the twentieth century was given by Gajewski and Życzkowski (1988) and Krużelecki (2004). A bibliography on the problems of topology and shape optimization of structures using FEM and BEM for 1999-2001 was collected by Mackerle (2003). Optimal design criteria for shapes of thin-walled beams cross-sections under strength and local and global stability constraints was presented by Cardoso (2000). Karim and Adeli (1999) presented global optimum design of cold-formed steel hat-shape beams under uniformly distributed load using a neural network model. Variational and parametric design of an open crosssection of a thin-walled beam under stability constraints was described by Magnucki and Magnucka-Blandzi (1999), Magnucki and Monczak (2000). Vinot et al. (2001) presented a methodology for optimizing the shape of thin-walled structures. Magnucki (2002) studied optimization of an open cross-section of a thin-walled beam with flat web and circular flange analytically and numerically. Knowledge-based global optimization of cold-formed steel columns under pure axial compression was presented by Liu et al. (2004). In result of the study, five anti-symmetrical open cross-sections were proposed. Theoretical and experimental study on the minimum weight of cold-formed channel thin-walled beams with and without lips were analysed by Tian and Lu (2004). Optimum design of cold-formed steel channel beams under uniformly distributed load
using micro Genetic Algorithm was presented by Lee et al. (2005). Global optimization of cold-formed steel thin-walled beams with lipped channel sections were described by Tran and Li (2006). Optimal design of open cross-sections of cold-formed thin-walled beams with respect to the dimensionless objective function as the quality measure was presented by Magnucka-Blandzi and Magnucki (2004b), Magnucki and Ostwald (2005a,b), Magnucki et al. (2006a,b), Magnucki and Paczos (2008). Kasperska et al. (2007), Ostwald et al. (2007), Ostwald and Magnucki (2008), Manevich and Raksha (2007) described bicriterial optimal design of open cross-sections of cold-formed beams. Strength, global and local buckling and optimization problems of cold-formed thin-walled beams with open cross-sections were collected and described by Magnucki and Ostwald (2005a,b), Ostwald and Magnucki (2008).

The present paper provides further development of optimal shaping of anti-symmetrical open cross-sections of cold-formed thin-walled beams in pure bending state. These beams of the length $L$, depth $H$, and wall thickness $t$ are simply supported and carry two equal moments $M_{0}$ applied to the beam ends (Fig. 1). The optimization includes three anti-symmetrical cross-sections: Z-section, S-section and clothoid-section.


Fig. 1. A scheme of the thin-walled beam

## 2. Geometric properties of three cross-sections

### 2.1. Anti-symmetrical Z-section

A scheme of the cross-section with principal axes $y z$ is shown in Fig. 2. The middle line of the Z-section is a broken line situated symmetrically with respect to the origin $O(0,0)$.

Geometric properties of the cross-section are defined by the following dimensionless parameters

$$
\begin{equation*}
x_{1}=\frac{b}{a} \quad x_{2}=\frac{c}{b} \quad x_{3}=\frac{t}{b} \quad x_{4}=\frac{d}{a} \tag{2.1}
\end{equation*}
$$

where: $a, b, c, d$-sizes of the cross-section, $t$ - thickness of the wall.


Fig. 2. A scheme of the Z-section

Depth of the beam is

$$
\begin{equation*}
H=2 a+t=a\left(2+x_{1} x_{3}\right) \tag{2.2}
\end{equation*}
$$

Total area and geometric stiffness for Saint-Venant torsion of the cross-section

$$
\begin{equation*}
A=2 a t f_{0}\left(x_{i}\right) \quad J_{S-V}=\frac{2}{3} a t^{3} f_{0}\left(x_{i}\right) \tag{2.3}
\end{equation*}
$$

where

$$
f_{0}\left(x_{i}\right)=x_{1}\left(1+x_{2}\right)+x_{4}+\sqrt{\left(1-x_{4}\right)^{2}+\frac{1}{4} x_{1}^{2}}
$$

The product of inertia with respect to the principal axes $y z$ is zero

$$
\begin{equation*}
J_{y z}=2 a^{3} t x_{1}\left[-3\left(x_{1} x_{2}-x_{4}\right)\left(2-x_{1} x_{2}-x_{4}\right)+\left(1-x_{1}\right) \sqrt{x_{1}^{2}+4\left(1-x_{4}\right)^{2}}\right]=0 \tag{2.4}
\end{equation*}
$$

from which

$$
\begin{equation*}
x_{2}=\frac{1}{x_{1}}\left(1+\sqrt{1-C_{0}}\right) \tag{2.5}
\end{equation*}
$$

where

$$
C_{0}=\left(2-x_{4}\right) x_{4}+\frac{1}{3}\left(1-x_{1}\right) \sqrt{x_{1}^{2}+4\left(1-x_{4}\right)^{2}}
$$

Moments of inertia of the plane area (Fig. 2) with respect to the $y$ and $z$ axes are

$$
\begin{equation*}
J_{y}=2 a^{3} t f_{2}\left(x_{i}\right) \quad J_{z}=2 a^{3} t f_{3}\left(x_{i}\right) \tag{2.6}
\end{equation*}
$$

where

$$
\begin{aligned}
f_{2}\left(x_{i}\right) & =\frac{1}{4} x_{1}^{2}\left[x_{1}\left(\frac{1}{3}+x_{2}\right)+x_{4}+\frac{1}{6} \sqrt{x_{1}^{2}+4\left(1-x_{4}\right)^{2}}\right] \\
f_{3}\left(x_{i}\right) & =x_{1}+\frac{1}{3}\left[2-\left(1-x_{1} x_{2}\right)^{3}-\left(1-x_{4}\right)^{3}\right]+\frac{1}{6}\left(1-x_{4}\right)^{2} \sqrt{x_{1}^{2}+4\left(1-x_{4}\right)^{2}}
\end{aligned}
$$

The warping function $\omega(s)$ for the Z-section (half section) is shown in Fig. 3.


Fig. 3. Geometric interpretation of the warping function $\omega(s)$
The warping function in characteristic points of the Z-section have the following values

$$
\begin{equation*}
\omega_{1}=0 \quad \omega_{i}=a^{i} \widetilde{\omega}_{i} \quad i=2,3,4 \tag{2.7}
\end{equation*}
$$

where

$$
\widetilde{\omega}_{2}=\frac{1}{2} x_{1} x_{4} \quad \widetilde{\omega}_{3}=\left(1+\frac{1}{2} x_{4}\right) x_{1} \quad \widetilde{\omega}_{4}=\left[1+\frac{1}{2}\left(x_{1} x_{2}+x_{4}\right)\right] x_{1}
$$

The warping moment of inertia

$$
\begin{equation*}
J_{\omega}=2 a^{5} t f_{5}\left(x_{i}\right) \tag{2.8}
\end{equation*}
$$

where

$$
f_{5}\left(x_{i}\right)=\frac{1}{3}\left[x_{4} \widetilde{\omega}_{2}^{2}+x_{1}\left(\widetilde{\omega}_{2}^{2}+\widetilde{\omega}_{2} \widetilde{\omega}_{3}+\widetilde{\omega}_{3}^{2}\right)+x_{1} x_{2}\left(\widetilde{\omega}_{3}^{2}+\widetilde{\omega}_{3} \widetilde{\omega}_{4}+\widetilde{\omega}_{4}^{2}\right)\right]
$$

The centroid and the shear center of the plane area of anti-symmetrical crosssections are located in the origin $O(0,0)$.

### 2.2. Anti-symmetrical S-section

A scheme of the cross-section with auxiliary axes $y_{1} z_{1}$ and principal axes $y z$ is shown in Fig. 4. The middle line of the S-section is a composite curve (two circles and one line segment) situated symmetrically with respect to the origin $O(0,0)$.


Fig. 4. A scheme of the S-section

Geometric properties of the cross-section are defined by the following dimensionless parameters

$$
\begin{equation*}
x_{1}=\frac{r}{a} \quad x_{3}=\frac{t}{r} \quad \text { and } \quad \beta \tag{2.9}
\end{equation*}
$$

where: $a, r$ - sizes of the cross-section, $\beta$ - angle, $t$ - thickness of the wall.
Depth of the beam is

$$
\begin{equation*}
H=2 a \cos \theta_{p}+2 r+t=2 a\left[\cos \theta_{p}+x_{1}\left(1+\frac{1}{2} x_{3}\right)\right] \tag{2.10}
\end{equation*}
$$

Total area and geometric stiffness for Saint-Venant torsion of the cross-section

$$
\begin{equation*}
A=2 a t f_{0}\left(x_{i}\right) \quad J_{S-V}=\frac{2}{3} a t^{3} f_{0}\left(x_{i}\right) \tag{2.11}
\end{equation*}
$$

where

$$
f_{0}\left(x_{i}\right)=\sqrt{1-x_{1}^{2}}+x_{1}(\pi+\beta-\alpha) \quad \cos \alpha=x_{1}
$$

The product of inertia with respect to the auxiliary axes $y_{1} z_{1}$

$$
\begin{equation*}
J_{y_{1} z_{1}}=2 a^{3} t f_{1}\left(x_{i}\right) \tag{2.12}
\end{equation*}
$$

where

$$
f_{1}\left(x_{i}\right)=\left\{\frac{1}{3}\left(1-x_{1}^{2}\right)^{2}+x_{1}\left[\cos \beta+x_{1}+\frac{1}{4} x_{1}\left(1+\cos 2 \beta-2 x_{1}^{2}\right)\right]\right\} x_{1}
$$

Moments of inertia of the plane area (Fig. 4) with respect to the $y_{1}$ and $z_{1}$ auxiliary axes are

$$
\begin{equation*}
J_{y_{1}}=2 a^{3} t f_{2}\left(x_{i}\right) \quad J_{z_{1}}=2 a^{3} t f_{3}\left(x_{i}\right) \tag{2.13}
\end{equation*}
$$

where

$$
\begin{aligned}
& f_{2}\left(x_{i}\right)=\left\{\frac{1}{3} \sqrt{\left(1-x_{1}^{2}\right)^{3}}+\frac{1}{4} x_{1}\left[2(\pi+\beta-\alpha)+2 x_{1} \sqrt{1-x_{1}^{2}}-\sin 2 \beta\right]\right\} x_{1}^{2} \\
& f_{3}\left(x_{i}\right)
\end{aligned}=\frac{1}{3} \sqrt{\left(1-x_{1}^{2}\right)^{5}}+\quad \begin{aligned}
+\frac{1}{2} x_{1}\left\{(\pi+\beta-\alpha)\left(2+x_{1}^{2}\right)+x_{1}\left[\sqrt{1-x_{1}^{2}}\left(4-x_{1}^{2}\right)+\left(4+x_{1} \cos \beta\right) \sin \beta\right]\right\}
\end{aligned}
$$

The angle $\theta_{p}$ defining the principal axes is

$$
\begin{equation*}
\tan 2 \theta_{p}=-\frac{2 J_{y_{1} z_{1}}}{J_{z_{1}}-J_{y_{1}}} \tag{2.14}
\end{equation*}
$$

and, principal moments of inertia

$$
\begin{align*}
J_{y} & =\frac{1}{2}\left(J_{z 1}+J_{y 1}\right)-\sqrt{\frac{1}{4}\left(J_{z 1}+J_{y 1}\right)^{2}+J_{y_{1} z_{1}}^{2}}  \tag{2.15}\\
J_{z} & =\frac{1}{2}\left(J_{z 1}+J_{y 1}\right)+\sqrt{\frac{1}{4}\left(J_{z 1}+J_{y 1}\right)^{2}+J_{y_{1} z_{1}}^{2}}
\end{align*}
$$

The warping function $\omega(\varphi)$ for the S-section (half section) is shown in Fig. 5.


Fig. 5. Geometric interpretation of the warping function $\omega(\varphi)$
The warping function for the S -section is defined as follows

$$
\begin{equation*}
\omega(\varphi)=\left[\sin \alpha-\sin (\alpha+\varphi)+x_{1} \varphi\right] x_{1} a^{2} \tag{2.16}
\end{equation*}
$$

The warping moment of inertia

$$
\begin{equation*}
J_{\omega}=2 a^{5} t f_{5}\left(x_{i}\right) \tag{2.17}
\end{equation*}
$$

where

$$
\begin{aligned}
& f_{5}\left(x_{i}\right)=\left(f_{51}-f_{52}+f_{53}+f_{54}\right) x_{1}^{3} \\
& f_{51}=\frac{1}{2}\left(x_{1} \sqrt{1-x_{1}^{2}}-\frac{1}{2} \sin 2 \beta\right) \\
& f_{52}=2\left[(\pi+\beta-\alpha) x_{1}+\sqrt{1-x_{1}^{2}}\right] \cos \beta \\
& f_{53}=\left[2 \sin \beta-(\pi+\beta-\alpha) \sqrt{1-x_{1}^{2}}\right] x_{1} \\
& f_{54}=\frac{1}{6}(\pi+\beta-\alpha)\left\{9+2\left[(\pi+\beta-\alpha)^{2}-3\right] x_{1}^{2}\right\}
\end{aligned}
$$

### 2.3. Anti-symmetrical Clothoid-section

A scheme of the cross-section with auxiliary axes $y_{1} z_{1}$ and principal axes $y z$ is shown in Fig. 6. The middle line of the Clothoid-section is a curve situated symmetrically with respect to the origin $O(0,0)$.


Fig. 6. A scheme of the Clothoid-section
In Cartesian auxiliary coordinates, the curve is parametrized as follows

$$
\begin{equation*}
y_{1}=a \sqrt{\pi} \int_{0}^{u_{1}} \sin \frac{\pi u^{2}}{2} d u \quad z_{1}=a \sqrt{\pi} \int_{0}^{u_{1}} \cos \frac{\pi u^{2}}{2} d u \tag{2.18}
\end{equation*}
$$

where $a$ is the scale parameter determining the outer size of the curve, $u$ - dimensionless parameter $\left(0 \leqslant u \leqslant u_{1}\right)$.

The principal curvature radius

$$
\begin{equation*}
R=\frac{a^{2}}{s} \tag{2.19}
\end{equation*}
$$

where arc length $s=a \sqrt{\pi} u_{1}$.
Geometric properties of the cross-section are defined by the following dimensionless parameters

$$
\begin{equation*}
x_{1}=u_{1} \quad x_{3}=\frac{t}{a} \tag{2.20}
\end{equation*}
$$

Depth of the beam is

$$
\begin{equation*}
H=2 d+t \tag{2.21}
\end{equation*}
$$

where: $u_{1}$ is the upper integration limit, as in (2.18), deciding on the "depth" of convolutions of the curve, $t$ - wall thickness.

The total area of the clothoid cross-section is

$$
\begin{equation*}
A=2 \int_{A} d A=2 t \int_{O P} d s=2 \sqrt{\pi} a t u_{1} \tag{2.22}
\end{equation*}
$$

The moments of inertia of the plane area with respect to the $z_{1}$ and $y_{1}$ axes are

$$
\begin{equation*}
I_{z 1}=\int_{A} y_{1}^{2} d A=2 a t \sqrt{\pi} \int_{0}^{u_{1}} a^{2} \pi\left(\int_{0}^{u} \sin \frac{\pi v^{2}}{2} d v\right) d u=2 a^{3} t \sqrt{\pi^{3}} \int_{0}^{u_{1}}[s(u)]^{2} d u \tag{2.23}
\end{equation*}
$$

where

$$
\begin{equation*}
s(u)=\int_{0}^{u} \sin \frac{\pi u^{2}}{2} d u \tag{2.24}
\end{equation*}
$$

and

$$
\begin{equation*}
I_{y 1}=\int_{A} z_{1}^{2} d A=2 a^{3} t \sqrt{\pi^{3}} \int_{0}^{u_{1}}[c(u)]^{2} d u \tag{2.25}
\end{equation*}
$$

where

$$
\begin{equation*}
c(u)=\int_{0}^{u} \cos \frac{\pi u^{2}}{2} d u \tag{2.26}
\end{equation*}
$$

The principal axes and principal moments of inertia are defined by the same expressions as for the S -section, i.e. (2.14) and (2.15).

The warping function of the clothoid section (Fig. 7) takes the following form

$$
\begin{equation*}
\omega=2\left(\frac{1}{2} z_{p} y_{p}-\int_{0}^{z_{p}} y_{1} d z_{1}\right) \tag{2.27}
\end{equation*}
$$

where

$$
z_{p}=a \sqrt{\pi} \int_{0}^{u_{p}} \cos \frac{\pi u^{2}}{2} d u \quad y_{p}=a \sqrt{\pi} \int_{0}^{u_{p}} \sin \frac{\pi u^{2}}{2} d u
$$



Fig. 7. Geometric interpretation of the clothoid warping function $\omega(u)$
According to earlier definitions (2.23) and (2.25)

$$
\begin{equation*}
z_{p} y_{p}=\pi a^{2} c\left(u_{p}\right) s\left(u_{p}\right) \quad d z_{1}=a \sqrt{\pi} \cos \frac{\pi u^{2}}{2} d u \tag{2.28}
\end{equation*}
$$

Hence, the warping function, being a function of the parameter $u_{p}$, may be formulated as follows

$$
\begin{equation*}
\omega\left(u_{p}\right)=\pi a^{2}\left[c\left(u_{p}\right) s\left(u_{p}\right)-2 \int_{0}^{u_{p}} \cos \frac{\pi u^{2}}{2} s(u) d u\right] \tag{2.29}
\end{equation*}
$$

Finally, the warping moment of inertia of the clothoid section is calculated as follows

$$
\begin{equation*}
I_{\omega}=\int_{A} \omega^{2} d A=2 t \int_{0}^{u_{p}} \omega^{2}(u) d s \tag{2.30}
\end{equation*}
$$

## 3. Formulation of the optimization problem

### 3.1. Optimization criterion

The minimal mass and maximal safe load are usually a basic objective in structure designing. The optimization criterion according to the papers of Magnucka-Blandzi and Magnucki (2004a,b), Magnucki et al. (2006a,b), has been formulated in the following form

$$
\begin{equation*}
\max _{x_{i}}\left\{\Phi_{1}\left(x_{i}\right), \Phi_{2}\left(x_{i}\right), \Phi_{3}\left(x_{i}\right), \Phi_{4}\left(x_{i}\right)\right\}=\Phi_{\max } \tag{3.1}
\end{equation*}
$$

and the objective function

$$
\begin{equation*}
\Phi_{j}\left(x_{i}\right)=\frac{M_{j}}{E \sqrt{A^{3}}} \tag{3.2}
\end{equation*}
$$

where $M_{j}$ are the allowable moments defined from the strength condition ( $j=1$ ), lateral buckling condition $(j=2)$, local buckling condition of the flange ( $j=3$ ), and local buckling condition of the web.

### 3.2. Constraints

Strength and buckling are main problems in thin-walled structures designing. Lateral buckling strengths of a cold-formed Z-section beam was presented by Pi et al. (1999). Li (2004) described lateral-torsion buckling of the cold-formed Z-beam. The effects of warping stress on the lateral torsional buckling, and local and distortional buckling of cold-formed Z-beams were described by Chu et al. (2004, 2006). Stasiewicz et al. (2004) described local buckling of a bent flange of a thin-walled beam. Analytical and numerical analysis of the stress state and global elastic buckling of a thin-walled beam with a mono-symmetrical open cross-section was presented by Magnucki et al. (2004). Critical stresses for open cylindrical shells with free edges were calculated by Magnucka-Blandzi and Magnucki (2004b), Magnucki and Mackiewicz (2006) and Joniak et al. (2008). Ventsel and Krauthammer (2001) collected and described strength and buckling problems of thin plates and shells.

The space of feasible solutions for optimal shapes of cross-sections of thinwalled beams is restrained. The strength condition has the following form

$$
\begin{equation*}
M_{0} \leqslant M_{1} \quad M_{1}=2 \frac{J_{z}}{H} \sigma_{\text {all }} \tag{3.3}
\end{equation*}
$$

where $\sigma_{\text {all }}$ is the allowable stress.
The global stability condition (lateral buckling condition) for a simply supported beam in pure bending state has the following form

$$
\begin{equation*}
M_{0} \leqslant M_{2} \quad M_{2}=\frac{M_{C R}^{(\text {Globl })}}{c_{s 1}} \tag{3.4}
\end{equation*}
$$

where $c_{s 1}$ is the safety coefficient, and the lateral buckling moment for a simply supported thin-walled beam in pure bending state is (Magnucki and Ostwald, 2005a,b)

$$
\begin{equation*}
M_{C R}^{(G l o b l)}=\frac{\pi E}{L} \sqrt{\frac{J_{y} J_{S-V}}{2(1+\nu)}\left[1+2(1+\nu) \frac{\pi^{2}}{L^{2}} \frac{J_{\omega}}{J_{S-V}}\right]} \tag{3.5}
\end{equation*}
$$

The local stability conditions for the Z-beam are as follows:

- for the bent flange, according to Magnucki and Ostwald (2005a,b) and Stasiewicz et al. (2004)

$$
\begin{equation*}
\sigma_{\text {max }}^{(Z-\text { flange })} \leqslant \frac{\sigma_{C R}^{(Z-\text { flange })}}{c_{s 2}} \quad \sigma_{C R}^{(Z-\text { flange })}=\frac{1+x_{2}}{1+3 x_{2}} x_{3}^{2} G \tag{3.6}
\end{equation*}
$$

where $\sigma_{C R}^{(Z-\text { flange })}$ is the critical stress, $G=E /[2(1+\nu)]-$ shear modulus of elasticity, $E$ - Young's modulus, $\nu$ - Poisson's ratio, $c_{s 2}$ - safety coefficient.

Taking into account the classical theory of plates, the local stability condition for the bent flange may be written down as

$$
M_{0} \leqslant M_{3} \quad M_{3}=\frac{\sigma_{C R}^{(Z-\text { flange })}}{c_{s 2}} \frac{J_{z}}{a-e_{f}}=\frac{2 a^{2} t}{c_{s 2}} G \frac{1+x_{2}}{1+3 x_{2}} x_{3}^{2} \frac{f_{3}\left(x_{i}\right)}{1-\widetilde{e}_{f}}
$$

where $\widetilde{e}_{f}$ is the dimensionless parameter of the centroid location of the flange

$$
\widetilde{e}_{f}=\frac{x_{2}}{2\left(1+x_{2}\right)} f_{1}\left(x_{i}\right)
$$

- for the flat web according to Ventsel and Krauthammer (2001)

$$
\begin{equation*}
\sigma_{\max }^{(Z-w e b)} \leqslant \frac{\sigma_{C R}^{(Z-w e b)}}{c_{s 2}} \quad \sigma_{C R}^{(Z-w e b)}=\frac{2 \pi^{2}}{1-\nu^{2}} E \frac{\left(x_{1} x_{3}\right)^{2}}{x_{1}^{2}+4\left(1-x_{4}\right)^{2}} \tag{3.8}
\end{equation*}
$$

where $\sigma_{C R}^{(Z-w e b)}$ is the critical stress.
Taking into account the classical theory of beams, the local stability condition for the flat web may be put down as

$$
\begin{align*}
& M_{0} \leqslant M_{4} \\
& M_{4}=\frac{\sigma_{C R}^{(Z-w e b)}}{c_{s 2}} \frac{J_{z}}{a-d}=\frac{4 \pi^{2} a^{2} t}{c_{s 2}\left(1-\nu^{2}\right)} E \frac{\left(x_{1} x_{3}\right)^{2}}{x_{1}^{2}+4\left(1-x_{4}\right)^{2}} \frac{f_{3}\left(x_{i}\right)}{1-x_{4}} \tag{3.9}
\end{align*}
$$

The local stability conditions for the $\mathbf{S}$-beam take the following forms:

- for the open circular cylindrical shell, regarding the results of Magnucka Blandzi and Magnucki (2004a), Magnucki and Mackiewicz (2006), Joniak et al. (2008)

$$
\begin{equation*}
\sigma_{\max }^{(S-\text { shell })} \leqslant \frac{\sigma_{C R}^{(S-\text { shell })}}{C_{s 2}} \quad \sigma_{C R}^{(S-\text { shell })}=\alpha_{C} \frac{E}{12.7 \sqrt{3\left(1-\nu^{2}\right)}} x_{3} \tag{3.10}
\end{equation*}
$$

where $\sigma_{C R}^{(S-\text { shell })}$ is the critical stress and $\alpha_{C}$ - coefficient

$$
\alpha_{C}=1+0.8\left(\beta-\frac{\pi}{2}\right)^{4}
$$

Taking into account the classical theory of beams, the local stability condition for the circular cylindrical flange may be written down as

$$
\begin{equation*}
M_{0} \leqslant M_{3} \quad M_{3}=\frac{\sigma_{C R}^{(S-\text { shell })}}{c_{s 2}} \frac{J_{z}}{a}=\frac{2 a t^{2}}{12.7 c_{s 2} \sqrt{3\left(1-\nu^{2}\right)}} E \alpha_{C} \frac{f_{3}\left(x_{i}\right)}{x_{1}} \tag{3.11}
\end{equation*}
$$

- for the flat web Ventsel and Krauthammer (2001)

$$
\begin{equation*}
\sigma_{\max }^{(S-w e b)} \leqslant \frac{\sigma_{C R}^{(S-w e b)}}{c_{s 2}} \quad \sigma_{C R}^{(S-w e b)}=\frac{\pi^{2}}{2\left(1-\nu^{2}\right)} E \frac{\left(x_{1} x_{3}\right)^{2}}{1-x_{1}} \tag{3.12}
\end{equation*}
$$

where $\sigma_{C R}^{(S-w e b)}$ is the critical stress.

Taking into account the classical theory of beams, the local stability condition for the flat web is

$$
\begin{align*}
& M_{0} \leqslant M_{4} \\
& M_{4}=\frac{\sigma_{C R}^{(S-w e b)}}{c_{s 2}} \frac{J_{z}}{a-d}=\frac{4 \pi^{2} a^{2} t}{c_{s 2}\left(1-\nu^{2}\right)} E \frac{\left(x_{1} x_{3}\right)^{2}}{x_{1}^{2}+4\left(1-x_{4}\right)^{2}} \frac{f_{3}\left(x_{i}\right)}{1-x_{4}} \tag{3.13}
\end{align*}
$$

The local stability conditions for the Clothoid-beam are as follows:

- for the open cylindrical shell, according to Magnucka Blandzi and Magnucki (2004b), Magnucki and Mackiewicz (2006), Joniak et al. (2008)

$$
\begin{equation*}
\sigma_{\text {edge }}^{(C l-\text { shell })} \leqslant \frac{\sigma_{C R, \text { edge }}^{(C l-\text { shell })}}{c_{s 2}} \quad \sigma_{C R, \text { edge }}^{(C l-\text { shell })}=\frac{E}{12.7 \sqrt{3\left(1-\nu^{2}\right)}} \frac{t}{R_{\text {edge }}} \tag{3.14}
\end{equation*}
$$

where $\sigma_{C R, \text { edge }}^{(C l-\text { shell })}$ is the critical stress.
Taking into account the classical theory of beams, the local stability condition for the circular cylindrical flange assumes the form

$$
\begin{align*}
& M_{0} \leqslant M_{3}  \tag{3.15}\\
& M_{3}=\frac{\sigma_{C R, e d g e}^{(C l-\text { shell })}}{c_{s 2}} \frac{J_{z}}{d}=\frac{2 a t^{2}}{12.7 c_{s 2} \sqrt{3\left(1-\nu^{2}\right)}} E \alpha_{C} \frac{f_{3}\left(x_{i}\right)}{x_{1}}
\end{align*}
$$

- for the cylindrical shell

$$
\begin{equation*}
\sigma_{\text {local }}^{(C l-\text { shell })} \leqslant \frac{\sigma_{C R, l o c a l}^{(C l-\text { shell })}}{c_{s 2}} \quad \sigma_{C R, l o c a l}^{(C l-\text { shell })}=\frac{E}{\sqrt{3\left(1-\nu^{2}\right)}} \frac{t}{R(y)} \tag{3.16}
\end{equation*}
$$

where $\sigma_{C R, \text { local }}^{(C l-\text { shell })}$ is the critical stress.
Taking into account the classical theory of beams, the local stability condition for the flat web is

$$
\begin{equation*}
M_{0} \leqslant M_{4} \quad M_{4}=\frac{\sigma_{C R, l o c a l}^{(C l-w e b)}}{c_{s 2}} \frac{J_{z}}{y(u)} \tag{3.17}
\end{equation*}
$$

## 4. Numerical solution of the optimization problem

Optimization of three anti-symmetrical open cross-sections has been performed for a family of cold-formed thin-walled beams: $\sigma_{\text {all }} / E=0.0015$, $\nu=0.3, c_{s 1}=1.5, c_{s 2}=2.1$, with relative lengths $\lambda=L / H=7.5,10.0$, $12.5,15.0,17.5,20.0$. The results of numerical calculations for the Z-beam are specified in Table 1, for the S-beam in Table 2, and for the Clothoid-beam in Table 3.

Table 1. Optimal parameters for the Z-beam

| $\lambda$ | 7.5 | 10.0 | 12.5 | 15.0 | 17.5 | 20.0 |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| $x_{1, \text { opt }}$ | 0.3897 | 0.5262 | 0.6607 | 0.7908 | 0.9161 | 1.0359 |
| $x_{2, \text { opt }}$ | 0.7259 | 0.4581 | 0.3051 | 0.2077 | 0.1400 | 0.0903 |
| $x_{3, \text { opt }}$ | 0.1175 | 0.1131 | 0.1098 | 0.1072 | 0.1052 | 0.1036 |
| $x_{4, \text { opt }}$ | 0.0458 | 0.0595 | 0.0725 | 0.0850 | 0.0964 | 0.1072 |
| $\Phi_{\max }$ | 0.0020725 | 0.001825 | 0.001645 | 0.001507 | 0.0013940 | 0.001301 |

Table 2. Optimal parameters for the S-beam

| $\lambda$ | 7.5 | 10.0 | 11.15 | 12.5 | 15.0 | 17.5 | 20.0 |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| $x_{1, \text { opt }}$ | 0.3696 | 0.3696 | 0.3696 | 0.4145 | 0.4971 | 0.5789 | 0.6598 |
| $\beta_{\text {opt }}$ | $\pi$ | $\pi$ | $\pi$ | $\pi$ | $\pi$ | $\pi$ | $\pi$ |
| $x_{3, \text { opt }}$ | 0.047951 | 0.047951 | 0.047951 | 0.04641 | 0.04383 | 0.04155 | 0.0395 |
| $\Phi_{\text {max }}$ | 0.003049 | 0.003049 | 0.003049 | 0.002857 | 0.002573 | 0.002356 | 0.002187 |

Table 3. Optimal parameters for the Clothoid-beam

| $\lambda$ | 7.5 | 10.0 | 12.5 | 15.0 | 17.5 | 20.0 |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| $x_{1, \text { opt }}$ | 2.6 | 2.6 | 2.6 | 2.6 | 2.6 | 2.6 |
| $x_{3, \text { opt }}$ | 0.014353 | 0.014353 | 0.014353 | 0.014353 | 0.014353 | 0.014353 |
| $\Phi_{\max }$ | 0.003553 | 0.003553 | 0.003553 | 0.003553 | 0.003553 | 0.003553 |

## 5. Conclusions

The criterion of effective shaping (optimal design) with dimensionless objective functions (26) enables sorting and comparing beams with arbitrary crosssections. This criterion is a quality measure of the cross-sections of beams.


Fig. 8. Dimensionless objective function $\Phi_{\max }$ for three considered types of beams

According to the plots in Fig. 8, the following conclusions may be drawn:

- In the case of the Z-Section, the lateral buckling is decisive for the beam of relative length $7.5 \leqslant \lambda$.
- For the S-Section of the relative length $\lambda \leqslant 11.15$, the lateral buckling imposes no constraint - the condition remains inactive. It is active only for $\lambda>11.15$.
- In the case of the Clothoid-Section, the lateral buckling remains inactive within the whole considered range of the relative length $\lambda$.

The beams with the Clothoid-section are definitely better than those with Zor S-sections.

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## Optymalizacja otwartych antysymetrycznych przekrojów belek cienkościennych walcowanych na zimno

Streszczenie
W artykule rozważane są belki o przekrojach poprzecznych w kształcie Z-, S- oraz w kształcie klotoidy. Zamieszczono krótki przegląd zagadnień optymalnego projektowania belek cienkościennych o przekrojach otwartych. Opisano właściwości geometryczne trzech rozważanych przekrojów. Zapisano warunki wytrzymałości oraz lokalnej i ogólnej stateczności belek cienkościennych. Sformułowano kryterium optymalizacyjne z wykorzystaniem bezwymiarowej funkcji celu będącej miarą jakości przekroju. Wyniki numerycznych obliczeń optymalnych zarysów przekrojów poprzecznych przedstawiono w tablicach i na rysunkach.

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