OPTIMIZATION OF ANTI-SYMMETRICAL OPEN CROSS-SECTIONS OF COLD-FORMED THIN-WALLED BEAMS

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The paper deals with cold-formed thin-walled beams with the Z-, S- and Clothoid-section. A short survey of optimal designs of thin-walled beams with open cross-sections is given. Geometric properties of three crosssections are described. Strength, local and global buckling conditions for thin-walled beams are presented. The optimal design criterion with a dimensionless objective function as a quality measure is defined. Results of numerical calculations for optimal shapes of three cross-sections are presented in tables and figures.

 $Key\ words:$ thin-walled beam, open cross-section, global and local buckling, optimal design

Notations

a,b,c,d	_	dimensions of cross-sections
r	_	radius of the circular arc
t	_	thickness of the beam wall
u	_	dimensionless parameter of the clothoid
A	_	area of the cross-section
Н	_	depth of the beam
L	_	length of the beam
J_{S-V}	_	geometric stiffness for Saint-Venant torsion
J_y, J_z	_	inertia moments

J_{ω}	_	warping moment of inertia
M_0	_	loading moment
R	_	principal radius of the clothoid
α, β	—	angles of the S-section
λ	_	relative length of the beam
$ heta_p$	_	angle to the principal axes
ω	_	warping function
Φ_j	_	objective function

1. Introduction

Shapes of open cross-sections of contemporary cold-formed thin-walled beams are rather complicated. They are usually mono-symmetrical, although sometimes anti-symmetrical too. The main constraints in designing thin-walled structures are strength and stability conditions. The beginnings of the optimal design of thin-walled structures reach back to 1959. The first paper on optimal design of a thin-walled beam with an open cross-section (I-section) in pure bending state was presented by Krishnan and Shetty (1959). A complete survey of optimal design problems of structures for the second half of the twentieth century was given by Gajewski and Życzkowski (1988) and Krużelecki (2004). A bibliography on the problems of topology and shape optimization of structures using FEM and BEM for 1999-2001 was collected by Mackerle (2003). Optimal design criteria for shapes of thin-walled beams cross-sections under strength and local and global stability constraints was presented by Cardoso (2000). Karim and Adeli (1999) presented global optimum design of cold-formed steel hat-shape beams under uniformly distributed load using a neural network model. Variational and parametric design of an open crosssection of a thin-walled beam under stability constraints was described by Magnucki and Magnucka-Blandzi (1999), Magnucki and Monczak (2000). Vinot et al. (2001) presented a methodology for optimizing the shape of thin-walled structures. Magnucki (2002) studied optimization of an open cross-section of a thin-walled beam with flat web and circular flange analytically and numerically. Knowledge-based global optimization of cold-formed steel columns under pure axial compression was presented by Liu et al. (2004). In result of the study, five anti-symmetrical open cross-sections were proposed. Theoretical and experimental study on the minimum weight of cold-formed channel thin-walled beams with and without lips were analysed by Tian and Lu (2004). Optimum design of cold-formed steel channel beams under uniformly distributed load

using micro Genetic Algorithm was presented by Lee *et al.* (2005). Global optimization of cold-formed steel thin-walled beams with lipped channel sections were described by Tran and Li (2006). Optimal design of open cross-sections of cold-formed thin-walled beams with respect to the dimensionless objective function as the quality measure was presented by Magnucka-Blandzi and Magnucki (2004b), Magnucki and Ostwald (2005a,b), Magnucki *et al.* (2006a,b), Magnucki and Paczos (2008). Kasperska *et al.* (2007), Ostwald *et al.* (2007), Ostwald and Magnucki (2008), Manevich and Raksha (2007) described bicriterial optimal design of open cross-sections of cold-formed beams. Strength, global and local buckling and optimization problems of cold-formed thin-walled beams with open cross-sections were collected and described by Magnucki and Ostwald (2005a,b), Ostwald and Magnucki (2008).

The present paper provides further development of optimal shaping of anti-symmetrical open cross-sections of cold-formed thin-walled beams in pure bending state. These beams of the length L, depth H, and wall thickness tare simply supported and carry two equal moments M_0 applied to the beam ends (Fig. 1). The optimization includes three anti-symmetrical cross-sections: Z-section, S-section and clothoid-section.

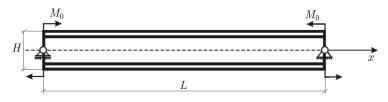


Fig. 1. A scheme of the thin-walled beam

2. Geometric properties of three cross-sections

2.1. Anti-symmetrical Z-section

A scheme of the cross-section with principal axes yz is shown in Fig. 2. The middle line of the Z-section is a broken line situated symmetrically with respect to the origin O(0,0).

Geometric properties of the cross-section are defined by the following dimensionless parameters

$$x_1 = \frac{b}{a}$$
 $x_2 = \frac{c}{b}$ $x_3 = \frac{t}{b}$ $x_4 = \frac{d}{a}$ (2.1)

where: a, b, c, d – sizes of the cross-section, t – thickness of the wall.

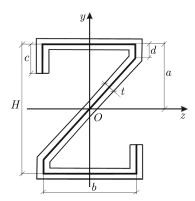


Fig. 2. A scheme of the Z-section

Depth of the beam is

$$H = 2a + t = a(2 + x_1 x_3) \tag{2.2}$$

Total area and geometric stiffness for Saint-Venant torsion of the cross-section

$$A = 2atf_0(x_i) \qquad J_{S-V} = \frac{2}{3}at^3 f_0(x_i) \qquad (2.3)$$

where

$$f_0(x_i) = x_1(1+x_2) + x_4 + \sqrt{(1-x_4)^2 + \frac{1}{4}x_1^2}$$

The product of inertia with respect to the principal axes yz is zero

$$J_{yz} = 2a^{3}tx_{1}\left[-3(x_{1}x_{2} - x_{4})(2 - x_{1}x_{2} - x_{4}) + (1 - x_{1})\sqrt{x_{1}^{2} + 4(1 - x_{4})^{2}}\right] = 0$$
(2.4)

from which

$$x_2 = \frac{1}{x_1} (1 + \sqrt{1 - C_0}) \tag{2.5}$$

where

$$C_0 = (2 - x_4)x_4 + \frac{1}{3}(1 - x_1)\sqrt{x_1^2 + 4(1 - x_4)^2}$$

Moments of inertia of the plane area (Fig. 2) with respect to the y and z axes are

$$J_y = 2a^3 t f_2(x_i) \qquad \qquad J_z = 2a^3 t f_3(x_i) \tag{2.6}$$

where

$$f_2(x_i) = \frac{1}{4}x_1^2 \left[x_1 \left(\frac{1}{3} + x_2 \right) + x_4 + \frac{1}{6}\sqrt{x_1^2 + 4(1 - x_4)^2} \right]$$

$$f_3(x_i) = x_1 + \frac{1}{3} \left[2 - (1 - x_1 x_2)^3 - (1 - x_4)^3 \right] + \frac{1}{6} (1 - x_4)^2 \sqrt{x_1^2 + 4(1 - x_4)^2}$$

The warping function $\omega(s)$ for the Z-section (half section) is shown in Fig. 3.

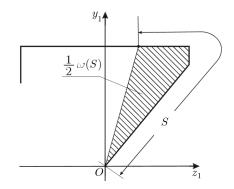


Fig. 3. Geometric interpretation of the warping function $\omega(s)$

The warping function in characteristic points of the Z-section have the following values

$$\omega_1 = 0 \qquad \qquad \omega_i = a^i \widetilde{\omega}_i \qquad \qquad i = 2, 3, 4 \tag{2.7}$$

where

$$\widetilde{\omega}_2 = \frac{1}{2}x_1x_4$$
 $\widetilde{\omega}_3 = \left(1 + \frac{1}{2}x_4\right)x_1$ $\widetilde{\omega}_4 = \left[1 + \frac{1}{2}(x_1x_2 + x_4)\right]x_1$

The warping moment of inertia

$$J_{\omega} = 2a^5 t f_5(x_i) \tag{2.8}$$

where

$$f_5(x_i) = \frac{1}{3} [x_4 \widetilde{\omega}_2^2 + x_1 (\widetilde{\omega}_2^2 + \widetilde{\omega}_2 \widetilde{\omega}_3 + \widetilde{\omega}_3^2) + x_1 x_2 (\widetilde{\omega}_3^2 + \widetilde{\omega}_3 \widetilde{\omega}_4 + \widetilde{\omega}_4^2)]$$

The centroid and the shear center of the plane area of anti-symmetrical crosssections are located in the origin O(0,0).

2.2. Anti-symmetrical S-section

A scheme of the cross-section with auxiliary axes y_1z_1 and principal axes yz is shown in Fig. 4. The middle line of the S-section is a composite curve (two circles and one line segment) situated symmetrically with respect to the origin O(0,0).

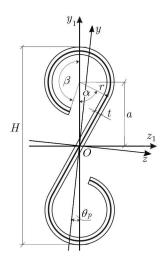


Fig. 4. A scheme of the S-section

Geometric properties of the cross-section are defined by the following dimensionless parameters

$$x_1 = \frac{r}{a}$$
 $x_3 = \frac{t}{r}$ and β (2.9)

where: a, r – sizes of the cross-section, β – angle, t – thickness of the wall. Depth of the beam is

$$H = 2a\cos\theta_p + 2r + t = 2a\left[\cos\theta_p + x_1\left(1 + \frac{1}{2}x_3\right)\right]$$
(2.10)

Total area and geometric stiffness for Saint-Venant torsion of the cross-section

$$A = 2atf_0(x_i) \qquad J_{S-V} = \frac{2}{3}at^3 f_0(x_i) \qquad (2.11)$$

where

$$f_0(x_i) = \sqrt{1 - x_1^2} + x_1(\pi + \beta - \alpha)$$
 $\cos \alpha = x_1$

The product of inertia with respect to the auxiliary axes $y_1 z_1$

$$J_{y_1 z_1} = 2a^3 t f_1(x_i) \tag{2.12}$$

where

$$f_1(x_i) = \left\{\frac{1}{3}(1-x_1^2)^2 + x_1\left[\cos\beta + x_1 + \frac{1}{4}x_1(1+\cos 2\beta - 2x_1^2)\right]\right\}x_1$$

Moments of inertia of the plane area (Fig. 4) with respect to the y_1 and z_1 auxiliary axes are

$$J_{y_1} = 2a^3 t f_2(x_i) \qquad \qquad J_{z_1} = 2a^3 t f_3(x_i) \qquad (2.13)$$

where

$$f_{2}(x_{i}) = \left\{\frac{1}{3}\sqrt{(1-x_{1}^{2})^{3}} + \frac{1}{4}x_{1}[2(\pi+\beta-\alpha)+2x_{1}\sqrt{1-x_{1}^{2}}-\sin 2\beta]\right\}x_{1}^{2}$$

$$f_{3}(x_{i}) = \frac{1}{3}\sqrt{(1-x_{1}^{2})^{5}} + \frac{1}{2}x_{1}\{(\pi+\beta-\alpha)(2+x_{1}^{2})+x_{1}[\sqrt{1-x_{1}^{2}}(4-x_{1}^{2})+(4+x_{1}\cos\beta)\sin\beta]\}$$

The angle θ_p defining the principal axes is

$$\tan 2\theta_p = -\frac{2J_{y_1 z_1}}{J_{z_1} - J_{y_1}} \tag{2.14}$$

and, principal moments of inertia

$$J_{y} = \frac{1}{2}(J_{z1} + J_{y1}) - \sqrt{\frac{1}{4}(J_{z1} + J_{y1})^{2} + J_{y_{1}z_{1}}^{2}}$$

$$J_{z} = \frac{1}{2}(J_{z1} + J_{y1}) + \sqrt{\frac{1}{4}(J_{z1} + J_{y1})^{2} + J_{y_{1}z_{1}}^{2}}$$
(2.15)

The warping function $\omega(\varphi)$ for the S-section (half section) is shown in Fig. 5.

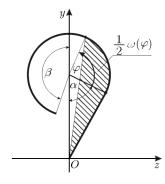


Fig. 5. Geometric interpretation of the warping function $\omega(\varphi)$

The warping function for the S-section is defined as follows

$$\omega(\varphi) = [\sin \alpha - \sin(\alpha + \varphi) + x_1 \varphi] x_1 a^2$$
(2.16)

The warping moment of inertia

$$J_{\omega} = 2a^5 t f_5(x_i) \tag{2.17}$$

where

$$f_{5}(x_{i}) = (f_{51} - f_{52} + f_{53} + f_{54})x_{1}^{3}$$

$$f_{51} = \frac{1}{2} \left(x_{1}\sqrt{1 - x_{1}^{2}} - \frac{1}{2}\sin 2\beta \right)$$

$$f_{52} = 2[(\pi + \beta - \alpha)x_{1} + \sqrt{1 - x_{1}^{2}}]\cos \beta$$

$$f_{53} = [2\sin\beta - (\pi + \beta - \alpha)\sqrt{1 - x_{1}^{2}}]x_{1}$$

$$f_{54} = \frac{1}{6}(\pi + \beta - \alpha)\{9 + 2[(\pi + \beta - \alpha)^{2} - 3]x_{1}^{2}\}$$

2.3. Anti-symmetrical Clothoid-section

A scheme of the cross-section with auxiliary axes y_1z_1 and principal axes yz is shown in Fig. 6. The middle line of the Clothoid-section is a curve situated symmetrically with respect to the origin O(0,0).

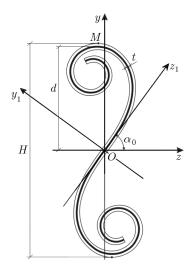


Fig. 6. A scheme of the Clothoid-section

In Cartesian auxiliary coordinates, the curve is parametrized as follows

$$y_1 = a\sqrt{\pi} \int_0^{u_1} \sin\frac{\pi u^2}{2} \, du \qquad z_1 = a\sqrt{\pi} \int_0^{u_1} \cos\frac{\pi u^2}{2} \, du \qquad (2.18)$$

where a is the scale parameter determining the outer size of the curve, u – dimensionless parameter $(0 \le u \le u_1)$.

The principal curvature radius

$$R = \frac{a^2}{s} \tag{2.19}$$

where arc length $s = a\sqrt{\pi} u_1$.

Geometric properties of the cross-section are defined by the following dimensionless parameters

$$x_1 = u_1$$
 $x_3 = \frac{t}{a}$ (2.20)

Depth of the beam is

$$H = 2d + t \tag{2.21}$$

where: u_1 is the upper integration limit, as in (2.18), deciding on the "depth" of convolutions of the curve, t – wall thickness.

The total area of the clothoid cross-section is

$$A = 2 \int_{A} dA = 2t \int_{OP} ds = 2\sqrt{\pi} atu_1 \qquad (2.22)$$

The moments of inertia of the plane area with respect to the z_1 and y_1 axes are

$$I_{z1} = \int_{A} y_1^2 \, dA = 2at\sqrt{\pi} \int_{0}^{u_1} a^2 \pi \left(\int_{0}^{u} \sin \frac{\pi v^2}{2} \, dv \right) \, du = 2a^3 t \sqrt{\pi^3} \int_{0}^{u_1} [s(u)]^2 \, du$$
(2.23)

where

$$s(u) = \int_{0}^{u} \sin \frac{\pi u^2}{2} \, du \tag{2.24}$$

and

$$I_{y1} = \int_{A} z_1^2 \, dA = 2a^3 t \sqrt{\pi^3} \int_{0}^{u_1} [c(u)]^2 \, du \tag{2.25}$$

where

$$c(u) = \int_{0}^{u} \cos \frac{\pi u^2}{2} \, du \tag{2.26}$$

The principal axes and principal moments of inertia are defined by the same expressions as for the S-section, i.e. (2.14) and (2.15).

The warping function of the clothoid section (Fig. 7) takes the following form

$$\omega = 2\left(\frac{1}{2}z_p y_p - \int_0^{z_p} y_1 \, dz_1\right) \tag{2.27}$$

where

$$z_p = a\sqrt{\pi} \int_{0}^{u_p} \cos\frac{\pi u^2}{2} \, du \qquad \qquad y_p = a\sqrt{\pi} \int_{0}^{u_p} \sin\frac{\pi u^2}{2} \, du$$

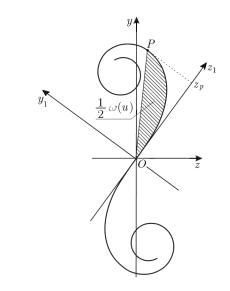


Fig. 7. Geometric interpretation of the clothoid warping function $\omega(u)$

According to earlier definitions (2.23) and (2.25)

$$z_p y_p = \pi a^2 c(u_p) s(u_p)$$
 $dz_1 = a \sqrt{\pi} \cos \frac{\pi u^2}{2} du$ (2.28)

Hence, the warping function, being a function of the parameter u_p , may be formulated as follows

$$\omega(u_p) = \pi a^2 \Big[c(u_p) s(u_p) - 2 \int_0^{u_p} \cos \frac{\pi u^2}{2} s(u) \, du \Big]$$
(2.29)

Finally, the warping moment of inertia of the clothoid section is calculated as follows

$$I_{\omega} = \int_{A} \omega^2 \, dA = 2t \int_{0}^{s_p} \omega^2(u) \, ds \tag{2.30}$$

3. Formulation of the optimization problem

3.1. Optimization criterion

The minimal mass and maximal safe load are usually a basic objective in structure designing. The optimization criterion according to the papers of Magnucka-Blandzi and Magnucki (2004a,b), Magnucki *et al.* (2006a,b), has been formulated in the following form

$$\max_{x_i} \{ \Phi_1(x_i), \Phi_2(x_i), \Phi_3(x_i), \Phi_4(x_i) \} = \Phi_{max}$$
(3.1)

and the objective function

$$\Phi_j(x_i) = \frac{M_j}{E\sqrt{A^3}} \tag{3.2}$$

where M_j are the allowable moments defined from the strength condition (j = 1), lateral buckling condition (j = 2), local buckling condition of the flange (j = 3), and local buckling condition of the web.

3.2. Constraints

Strength and buckling are main problems in thin-walled structures designing. Lateral buckling strengths of a cold-formed Z-section beam was presented by Pi *et al.* (1999). Li (2004) described lateral-torsion buckling of the cold-formed Z-beam. The effects of warping stress on the lateral torsional buckling, and local and distortional buckling of cold-formed Z-beams were described by Chu *et al.* (2004, 2006). Stasiewicz *et al.* (2004) described local buckling of a bent flange of a thin-walled beam. Analytical and numerical analysis of the stress state and global elastic buckling of a thin-walled beam with a mono-symmetrical open cross-section was presented by Magnucki *et al.* (2004). Critical stresses for open cylindrical shells with free edges were calculated by Magnucka-Blandzi and Magnucki (2004b), Magnucki and Mackiewicz (2006) and Joniak *et al.* (2008). Ventsel and Krauthammer (2001) collected and described strength and buckling problems of thin plates and shells. The space of feasible solutions for optimal shapes of cross-sections of thinwalled beams is restrained. The strength condition has the following form

$$M_0 \leqslant M_1 \qquad \qquad M_1 = 2\frac{J_z}{H}\sigma_{all} \tag{3.3}$$

where σ_{all} is the allowable stress.

The global stability condition (lateral buckling condition) for a simply supported beam in pure bending state has the following form

$$M_0 \leqslant M_2 \qquad \qquad M_2 = \frac{M_{CR}^{(Globl)}}{c_{s1}} \tag{3.4}$$

where c_{s1} is the safety coefficient, and the lateral buckling moment for a simply supported thin-walled beam in pure bending state is (Magnucki and Ostwald, 2005a,b)

$$M_{CR}^{(Globl)} = \frac{\pi E}{L} \sqrt{\frac{J_y J_{S-V}}{2(1+\nu)}} \Big[1 + 2(1+\nu) \frac{\pi^2}{L^2} \frac{J_\omega}{J_{S-V}} \Big]$$
(3.5)

The local stability conditions for the **Z-beam** are as follows:

• for the bent flange, according to Magnucki and Ostwald (2005a,b) and Stasiewicz *et al.* (2004)

$$\sigma_{max}^{(Z-flange)} \leqslant \frac{\sigma_{CR}^{(Z-flange)}}{c_{s2}} \qquad \qquad \sigma_{CR}^{(Z-flange)} = \frac{1+x_2}{1+3x_2} x_3^2 G \qquad (3.6)$$

where $\sigma_{CR}^{(Z-flange)}$ is the critical stress, $G = E/[2(1+\nu)]$ – shear modulus of elasticity, E – Young's modulus, ν – Poisson's ratio, c_{s2} – safety coefficient.

Taking into account the classical theory of plates, the local stability condition for the bent flange may be written down as

$$M_0 \leqslant M_3 \qquad M_3 = \frac{\sigma_{CR}^{(Z-flange)}}{c_{s2}} \frac{J_z}{a - e_f} = \frac{2a^2t}{c_{s2}} G \frac{1 + x_2}{1 + 3x_2} x_3^2 \frac{f_3(x_i)}{1 - \tilde{e}_f}$$
(3.7)

where $\,\widetilde{e}_f$ is the dimensionless parameter of the centroid location of the flange

$$\tilde{e}_f = \frac{x_2}{2(1+x_2)} f_1(x_i)$$

• for the flat web according to Ventsel and Krauthammer (2001)

$$\sigma_{max}^{(Z-web)} \leqslant \frac{\sigma_{CR}^{(Z-web)}}{c_{s2}} \qquad \qquad \sigma_{CR}^{(Z-web)} = \frac{2\pi^2}{1-\nu^2} E \frac{(x_1 x_3)^2}{x_1^2 + 4(1-x_4)^2}$$
(3.8)

where $\sigma_{CR}^{(Z-web)}$ is the critical stress.

Taking into account the classical theory of beams, the local stability condition for the flat web may be put down as

$$M_{0} \leqslant M_{4}$$

$$M_{4} = \frac{\sigma_{CR}^{(Z-web)}}{c_{s2}} \frac{J_{z}}{a-d} = \frac{4\pi^{2}a^{2}t}{c_{s2}(1-\nu^{2})} E \frac{(x_{1}x_{3})^{2}}{x_{1}^{2}+4(1-x_{4})^{2}} \frac{f_{3}(x_{i})}{1-x_{4}}$$
(3.9)

The local stability conditions for the **S-beam** take the following forms:

• for the open circular cylindrical shell, regarding the results of Magnucka Blandzi and Magnucki (2004a), Magnucki and Mackiewicz (2006), Joniak *et al.* (2008)

$$\sigma_{max}^{(S-shell)} \leqslant \frac{\sigma_{CR}^{(S-shell)}}{C_{s2}} \qquad \sigma_{CR}^{(S-shell)} = \alpha_C \frac{E}{12.7\sqrt{3(1-\nu^2)}} x_3 \quad (3.10)$$

where $\sigma_{CR}^{(S-shell)}$ is the critical stress and α_C – coefficient

$$\alpha_C = 1 + 0.8 \left(\beta - \frac{\pi}{2}\right)^4$$

Taking into account the classical theory of beams, the local stability condition for the circular cylindrical flange may be written down as

$$M_0 \leqslant M_3 \qquad M_3 = \frac{\sigma_{CR}^{(S-shell)}}{c_{s2}} \frac{J_z}{a} = \frac{2at^2}{12.7c_{s2}\sqrt{3(1-\nu^2)}} E\alpha_C \frac{f_3(x_i)}{x_1}$$
(3.11)

• for the flat web Ventsel and Krauthammer (2001)

$$\sigma_{max}^{(S-web)} \leqslant \frac{\sigma_{CR}^{(S-web)}}{c_{s2}} \qquad \qquad \sigma_{CR}^{(S-web)} = \frac{\pi^2}{2(1-\nu^2)} E \frac{(x_1 x_3)^2}{1-x_1} \qquad (3.12)$$

where $\sigma_{CR}^{(S-web)}$ is the critical stress.

Taking into account the classical theory of beams, the local stability condition for the flat web is

$$M_{0} \leqslant M_{4}$$

$$M_{4} = \frac{\sigma_{CR}^{(S-web)}}{c_{s2}} \frac{J_{z}}{a-d} = \frac{4\pi^{2}a^{2}t}{c_{s2}(1-\nu^{2})} E \frac{(x_{1}x_{3})^{2}}{x_{1}^{2}+4(1-x_{4})^{2}} \frac{f_{3}(x_{i})}{1-x_{4}}$$
(3.13)

The local stability conditions for the **Clothoid-beam** are as follows:

• for the open cylindrical shell, according to Magnucka Blandzi and Magnucki (2004b), Magnucki and Mackiewicz (2006), Joniak *et al.* (2008)

$$\sigma_{edge}^{(Cl-shell)} \leqslant \frac{\sigma_{CR,edge}^{(Cl-shell)}}{c_{s2}} \qquad \sigma_{CR,edge}^{(Cl-shell)} = \frac{E}{12.7\sqrt{3(1-\nu^2)}} \frac{t}{R_{edge}} \tag{3.14}$$

where $\sigma_{CR,edge}^{(Cl-shell)}$ is the critical stress.

Taking into account the classical theory of beams, the local stability condition for the circular cylindrical flange assumes the form

$$M_{0} \leq M_{3}$$

$$M_{3} = \frac{\sigma_{CR,edge}^{(Cl-shell)}}{c_{s2}} \frac{J_{z}}{d} = \frac{2at^{2}}{12.7c_{s2}\sqrt{3(1-\nu^{2})}} E\alpha_{C} \frac{f_{3}(x_{i})}{x_{1}}$$
(3.15)

• for the cylindrical shell

$$\sigma_{local}^{(Cl-shell)} \leqslant \frac{\sigma_{CR,local}^{(Cl-shell)}}{c_{s2}} \qquad \qquad \sigma_{CR,local}^{(Cl-shell)} = \frac{E}{\sqrt{3(1-\nu^2)}} \frac{t}{R(y)} \quad (3.16)$$

where $\sigma_{CR,local}^{(Cl-shell)}$ is the critical stress.

Taking into account the classical theory of beams, the local stability condition for the flat web is

$$M_0 \leqslant M_4 \qquad \qquad M_4 = \frac{\sigma_{CR,local}^{(Cl-web)}}{c_{s2}} \frac{J_z}{y(u)} \tag{3.17}$$

4. Numerical solution of the optimization problem

Optimization of three anti-symmetrical open cross-sections has been performed for a family of cold-formed thin-walled beams: $\sigma_{all}/E = 0.0015$, $\nu = 0.3$, $c_{s1} = 1.5$, $c_{s2} = 2.1$, with relative lengths $\lambda = L/H = 7.5$, 10.0, 12.5, 15.0, 17.5, 20.0. The results of numerical calculations for the Z-beam are specified in Table 1, for the S-beam in Table 2, and for the Clothoid-beam in Table 3.

λ	7.5	10.0	12.5	15.0	17.5	20.0
$x_{1,opt}$	0.3897	0.5262	0.6607	0.7908	0.9161	1.0359
$x_{2,opt}$	0.7259	0.4581	0.3051	0.2077	0.1400	0.0903
$x_{3,opt}$	0.1175	0.1131	0.1098	0.1072	0.1052	0.1036
$x_{4,opt}$	0.0458	0.0595	0.0725	0.0850	0.0964	0.1072
Φ_{max}	0.0020725	0.001825	0.001645	0.001507	0.0013940	0.001301

Table 1. Optimal parameters for the Z-beam

Table 2.	Optimal	parameters	for	the	S-beam
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λ	7.5	10.0	11.15	12.5	15.0	17.5	20.0
$x_{1,opt}$	0.3696	0.3696	0.3696	0.4145	0.4971	0.5789	0.6598
β_{opt}	π						
$x_{3,opt}$	0.047951	0.047951	0.047951	0.04641	0.04383	0.04155	0.0395
Φ_{max}	0.003049	0.003049	0.003049	0.002857	0.002573	0.002356	0.002187

Table 3. Optimal parameters for the Clothoid-beam

λ	7.5	10.0	12.5	15.0	17.5	20.0
$x_{1,opt}$	2.6	2.6	2.6	2.6	2.6	2.6
$x_{3,opt}$	0.014353	0.014353	0.014353	0.014353	0.014353	0.014353
Φ_{max}	0.003553	0.003553	0.003553	0.003553	0.003553	0.003553

5. Conclusions

The criterion of effective shaping (optimal design) with dimensionless objective functions (26) enables sorting and comparing beams with arbitrary cross-sections. This criterion is a quality measure of the cross-sections of beams.

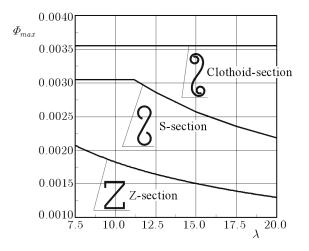


Fig. 8. Dimensionless objective function Φ_{max} for three considered types of beams

According to the plots in Fig. 8, the following conclusions may be drawn:

- In the case of the Z-Section, the lateral buckling is decisive for the beam of relative length $7.5 \leq \lambda$.
- For the S-Section of the relative length $\lambda \leq 11.15$, the lateral buckling imposes no constraint the condition remains inactive. It is active only for $\lambda > 11.15$.
- In the case of the Clothoid-Section, the lateral buckling remains inactive within the whole considered range of the relative length λ .

The beams with the Clothoid-section are definitely better than those with Zor S-sections.

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Optymalizacja otwartych antysymetrycznych przekrojów belek cienkościennych walcowanych na zimno

Streszczenie

W artykule rozważane są belki o przekrojach poprzecznych w kształcie Z-, S- oraz w kształcie klotoidy. Zamieszczono krótki przegląd zagadnień optymalnego projektowania belek cienkościennych o przekrojach otwartych. Opisano właściwości geometryczne trzech rozważanych przekrojów. Zapisano warunki wytrzymałości oraz lokalnej i ogólnej stateczności belek cienkościennych. Sformułowano kryterium optymalizacyjne z wykorzystaniem bezwymiarowej funkcji celu będącej miarą jakości przekroju. Wyniki numerycznych obliczeń optymalnych zarysów przekrojów poprzecznych przedstawiono w tablicach i na rysunkach.

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