

Exponential growth of solution for a couple of semi-linear pseudo-parabolic equations with memory and source terms

Amar Ouaoa  ¹ and Wissem Boughamsa ²

¹August 20, 1955 University, Skikda, Faculty of the sciences and technology, (LAMAHIS) Laboratory, Algeria

²August 20, 1955 University, Skikda, Faculty of the sciences, (LAMAHIS) Laboratory, Algeria

Received 30 January 2022, Accepted 21 April 2022, Published 18 May 2022

Abstract. This work is concerned with coupled semi-linear pseudo-parabolic equations with memory terms in both equations, associated with the homogeneous Dirichlet boundary condition. We show that the solution grows exponentially under specific conditions regarding the relaxation functions and initial energy. In order to prove the result, we use the energy method based on the construction of a suitable Lyapunov function. The most important behavior of the evolution system is the exponential growth phenomena because of its wide range of applications in modern science, such as chemistry, biology, ecology, and other areas of engineering and physical sciences.

Keywords: Memory terms, parabolic equation, exponential growth, source terms.

2020 Mathematics Subject Classification: 35B40, 35L90.


1 Introduction

We consider the following boundary value problem:

$$\begin{cases} u_t - \Delta u - \Delta u_t + \int_0^t g(t-s) \Delta u(s) ds + |u|^{m-2} u_t = f_1(u, v), & \text{in } \Omega \times (0, T) \\ v_t - \Delta v - \Delta v_t + \int_0^t h(t-s) \Delta v(s) ds + |v|^{k-2} v_t = f_2(u, v), & \text{in } \Omega \times (0, T) \\ u(x, t) = 0, v(x, t) = 0, & \text{in } \partial\Omega \times (0, T) \\ u(x, 0) = u_0, v(x, 0) = v_0, & \text{in } \Omega, \end{cases} \quad (1.1)$$

where Ω is a bounded domain in \mathbb{R}^n , $n \geq 1$ with smooth boundary $\partial\Omega$, m and k are real positive constants. The relaxation functions g and h satisfying some conditions we suppose later, and the two functions $f_1(u, v)$ and $f_2(u, v)$ are given by

$$\begin{cases} f_1(u, v) = |u + v|^{2(r+1)} (u + v) + |u|^r u |v|^{r+2}. \\ f_2(u, v) = |u + v|^{2(r+1)} (u + v) + |u|^{r+2} v |v|^r. \end{cases} \quad (1.2)$$

 Corresponding author. Email: ouaouaama21@gmail.com

Recently, Pişkin and Ekinçi in [15] treated the following system

$$\begin{cases} u_t - \Delta u + |u|^{q-2} u_t = f_1(u, v), \\ v_t - \Delta v + |v|^{q-2} v_t = f_2(u, v). \end{cases} \quad (1.3)$$

They proved the exponential growth of the solution with initial negative energy.

In the absence of $|u|^{q-2} u_t$ and $|v|^{q-2} v_t$ terms the system (1.3) becomes

$$\begin{cases} u_t - \Delta u = f_1(u, v), \\ v_t - \Delta v = f_2(u, v). \end{cases}$$

This type of equation is naturally found in physics, chemistry, biology, ecology, and other areas of engineering and physical sciences (see [4–6]). Korpusev in [9] has determined sufficient conditions for the blow-up of a finite time and the solvability of the following generalized Boussinesq equation

$$\frac{\partial}{\partial t} (\Delta u - u - |u|^p u) + u(u + \alpha)(u - \beta) = 0, \quad (1.4)$$

with initial boundary value, for $\alpha, \beta > 0$ in \mathbb{R}^3 , by concativity method [7, 10]. The result is extended by [20, 21]. Based on idea in [3] for the equation (1.5) the authors of [20] proved that the L^p -norm of the solution for (1.4) grows as an exponential function.

$$u_t - \Delta u - \Delta u_t + |u|^{m-2} u_t = |u|^{p-2} u. \quad (1.5)$$

In [16] Polat established a blow up result of the solution with vanishing initial energy in a bounded domain of \mathbb{R} with the absence of the term $-\Delta u_t$ in equation (1.5). Many authors [1, 2, 8, 17] have considered the following initial value problem for the generalized Boussinesq equation with nonlinear Newmann condition

$$\begin{cases} u_t - \Delta u - \Delta u_t + |u|^{m-2} u_t = f(u), \\ u = 0, \quad x \in \Gamma_0, \\ \frac{\partial u}{\partial \nu} + g(u) = 0, \quad x \in \Gamma_1, \\ u(x, 0) = u_0(x), \quad x \in \Omega. \end{cases} \quad (1.6)$$

In [2] the authors have established both the existence of the solution and a generalisation of the energy functions under some restrictions on the initial data. They have also proved a blow-up result. Ouaoua et al in [13] considered the following nonlinear Kirchhoff type reaction-diffusion equation

$$u_t - M \left(\int_{\Omega} |\nabla u|^2 dx \right) \Delta u + |u|^{m-2} u_t = |u|^{r-2} u, \quad (x, t) \in \Omega \times (0, T), \quad (1.7)$$

where $M(s) = a + bs^\gamma$, a, b and γ are positive constants. Under suitable assumptions on the initial data, they obtained global existence and stability of solutions with positive initial energy. In the case of the variable exponents, also, Ouaoua and Maouni in [14] considered the following equation

$$u_t - \operatorname{div} \left(|\nabla u|^{p(x)-2} \nabla u \right) + \omega |u|^{m(x)-2} u_t = b |u|^{r(x)-2} u \quad \text{in } \Omega \times (0, T). \quad (1.8)$$

They proved blow-up, exponential growth of solution with negative initial energy in the both case of equation. Messaoudi in [11] proved under suitable conditions on g and p , a blow-up result for solutions of the semilinear viscoelastic equation

$$u_t - \Delta u + \int_0^t g(t-s) \Delta u(s, x) dx = |u|^{p-2} u \text{ in } \Omega \times (0, T), \quad (1.9)$$

with positive initial energy.

The paper is organized as follows. In section 2, we present the necessary assumptions and lemmas, while in section 3, we prove the main result.

2 Preliminaries and assumptions

This section will give some notations and statements of assumptions for the relaxation functions. We denote $L^p(\Omega)$ by L^p and $H_0^1(\Omega)$ by H_0^1 . The norm and inner product of $L^p(\Omega)$ are denoted by $\|\cdot\|_p$ and $(a, b) = \int a(x)b(x) dx$ respectively.

For the relaxation functions g and h we assume that

(A₁) $g, h : \mathbb{R}^+ \rightarrow \mathbb{R}^+$ are of class C^1 and satisfying

$$\begin{cases} g(s) \geq 0, 1 - \int_0^\infty g(s) ds = l > 0, \\ h(s) \geq 0, 1 - \int_0^\infty h(s) ds = k > 0, \end{cases} \quad (2.1)$$

and

$$g'(s), h'(s) \leq 0, \forall s \geq 0. \quad (2.2)$$

(A₂)

$$\begin{cases} -1 < r \text{ if } n = 1, 2, \\ -1 < r \leq \frac{3-n}{n-2} \text{ if } n \geq 3. \end{cases} \quad (2.3)$$

Throughout this paper, we use the following notations

$$(\phi \circ \psi)(t) = \int_0^t \phi(t-s) \int_\Omega |\psi(t) - \psi(s)|^2 dx ds.$$

We can easily verify that

$$uf_1(u, v) + vf_2(u, v) = 2(r+2)F(u, v), \forall (u, v) \in \mathbb{R}^2,$$

where

$$F(u, v) = \frac{1}{2(r+2)} \left[a|u+v|^{2(r+2)} + 2b|uv|^{r+2} \right].$$

Lemma 2.1. [12]: *There exist two positive constants c_0 and c_1 such that*

$$\frac{c_0}{2(r+2)} \left(|u|^{2(r+2)} + |v|^{2(r+2)} \right) \leq F(u, v) \leq \frac{c_1}{2(r+2)} \left(|u|^{2(r+2)} + |v|^{2(r+2)} \right). \quad (2.4)$$

3 Main result and proof

In this section we state and prove our main result. First we state the definition for the strong solution of the problem (1.1).

Definition 3.1. A strong solution of the problem (1.1) on $\Omega \times [0, T]$, is a pair of (u, v) satisfying

$$(u, v) \in \left(C \left([0, T]; H_0^1 \right) \cap C^1 [0, T]; L^2 \right)^2, \left(|u|^{m-2} u_t, |v|^{k-2} v_t \right) \in [L^2(\Omega \times [0, T])]^2$$

$$(u(x, 0), v(x, 0)) = (u_0(x), v_0(x)),$$

$$\int_0^t \int_{\Omega} u_t \psi dx ds - \int_0^t \int_{\Omega} \int_0^s g(t-\tau) \nabla u(\tau) \nabla \psi(s) d\tau dx ds$$

$$+ \int_0^t \int_{\Omega} \nabla u \nabla \psi dx ds - \int_0^t \int_{\Omega} f_1(u, v) \psi dx ds + \int_0^t \int_{\Omega} |u|^{m-2} u_t \psi dx ds = 0.$$

and

$$\int_0^t \int_{\Omega} v_t \phi dx ds - \int_0^t \int_{\Omega} \int_0^s h(t-\tau) \nabla v(\tau) \nabla \phi(s) d\tau dx ds$$

$$+ \int_0^t \int_{\Omega} \nabla v \nabla \phi dx ds - \int_0^t \int_{\Omega} f_2(u, v) \phi dx ds + \int_0^t \int_{\Omega} |v|^{k-2} v_t \phi dx ds = 0.$$

for all $t \in [0, T]$ and all $\psi, \phi \in C([0, T], H_0^1)$.

The energy functional $E(t)$ associated with a solution (u, v) of (1.1) is given by

$$E(t) = \frac{1}{2} \left(1 - \int_0^t g(s) ds \right) \|\nabla u\|_2^2 + \frac{1}{2} \left(1 - \int_0^t h(s) ds \right) \|\nabla v\|_2^2 + \frac{1}{2} (g \circ \nabla u)$$

$$+ \frac{1}{2} (h \circ \nabla v) - \frac{1}{2(r+2)} \int_{\Omega} [u f_1(u, v) + v f_2(u, v)] dx. \quad (3.1)$$

Lemma 3.2. Let (u, v) be a solution of (1.1), then $E(t)$ is a nonincreasing function for $t > 0$ satisfying

$$E'(t) = \frac{1}{2} \left((g' \circ \nabla u)(t) + (h' \circ \nabla v)(t) - g(t) \|\nabla u(t)\|_2^2 - h(t) \|\nabla v(t)\|_2^2 \right)$$

$$- \|u_t\|_2^2 - \|v_t\|_2^2 - \int_{\Omega} |u|^{m-2} u_t^2 dx - \int_{\Omega} |v|^{k-2} v_t^2 dx \leq 0, \quad (3.2)$$

and

$$E(t) \leq E(0) \leq 0, \quad \forall t \geq 0.$$

Proof. Multiplying first equation of (1.1) by u_t and second by v_t , integrating over Ω and using integration by parts, we obtain

$$\frac{1}{2} \frac{d}{dt} \left[\left(1 - \int_0^t g(s) ds \right) \|\nabla u\|_2^2 + \left(1 - \int_0^t h(s) ds \right) \|\nabla v\|_2^2 \right.$$

$$- \frac{1}{2(r+2)} \int_{\Omega} [u f_1(u, v) + v f_2(u, v)] dx + \frac{1}{2} \int_{\Omega} \int_0^t g(t-s) |\nabla u(t) - \nabla u(s)|^2 ds dx$$

$$- \frac{1}{2} \int_{\Omega} \int_0^t g(t-s) |\nabla u(t)|^2 ds dx + \frac{1}{2} \int_{\Omega} \int_0^t h(t-s) |\nabla v(t) - \nabla v(s)|^2 ds dx$$

$$- \left. \frac{1}{2} \int_{\Omega} \int_0^t h(t-s) |\nabla v(t)|^2 ds dx \right]$$

$$- \frac{1}{2} \int_{\Omega} \int_0^t g'(t-s) |\nabla u(t) - \nabla u(s)|^2 ds dx + \frac{1}{2} \int_0^t g(t-s) ds \|\nabla u(t)\|_2^2$$

$$- \frac{1}{2} \int_{\Omega} \int_0^t h'(t-s) |\nabla v(t) - \nabla v(s)|^2 ds dx + \frac{1}{2} \int_0^t h(t-s) ds \|\nabla v(t)\|_2^2$$

$$- \int_{\Omega} |u|^{m-2} u_t^2 dx - \int_{\Omega} |v|^{k-2} v_t^2 dx - \|u_t\|_2^2 - \|v_t\|_2^2,$$

then by definition of $E(t)$, we get

$$\begin{aligned} E'(t) &= \frac{1}{2} \left((g' \circ \nabla u)(t) + (h' \circ \nabla v)(t) - g(t) \|\nabla u(t)\|_2^2 - h(t) \|\nabla v(t)\|_2^2 \right) \\ &\quad - \|u_t\|_2^2 - \|v_t\|_2^2 - \int_{\Omega} |u|^{m-2} u_t^2 dx - \int_{\Omega} |v|^{k-2} v_t^2 dx. \end{aligned} \quad (3.3)$$

This infer by (2.2) that $E'(t) \leq 0, \forall t \geq 0$ \square

Theorem 3.3. Suppose that assumptions (A_1) and (A_2) hold, $(u_0, v_0) \in (H_0^1)^2$ and (u, v) is a local strong solution of the system (1.1), and $E(0) < 0$.

Furthermore, we assume that $\max(\int_0^\infty g(s) ds, \int_0^\infty h(s) ds) < p/(p+1)$ and

$$2(r+2) > \max(m, k). \quad (3.4)$$

Then the solution of the system (1.1) exponentially grows.

Proof. We set

$$H(t) = -E(t). \quad (3.5)$$

By the definition of $H(t)$ and (3.5)

$$H(t) = -E(t) \geq 0. \quad (3.6)$$

By using (3.1) and (3.5) we get

$$H(t) - \frac{1}{2(r+2)} \int_{\Omega} [uf_1(u, v) + vf_2(u, v)] dx \leq 0. \quad (3.7)$$

So that

$$0 < H(0) \leq H(t) \leq \frac{1}{2(r+2)} \int_{\Omega} [uf_1(u, v) + vf_2(u, v)] dx. \quad (3.8)$$

We define a Lyapunov function

$$L(t) = H(t) + \frac{\varepsilon}{2} \left(\|u\|_2^2 + \|v\|_2^2 + \|\nabla u\|_2^2 + \|\nabla v\|_2^2 \right), \quad (3.9)$$

for ε small to be chosen later.

By taking the time derivative of (3.9) and (1.1), we obtain

$$\begin{aligned} L'(t) &= H'(t) + \varepsilon \int_{\Omega} u_t u dx + \varepsilon \int_{\Omega} v_t v dx + \varepsilon \int_{\Omega} \nabla u_t \nabla u dx + \varepsilon \int_{\Omega} \nabla v_t \nabla v dx \\ &\geq H'(t) - \varepsilon \left(\|\nabla u\|_2^2 + \|\nabla v\|_2^2 \right) + \varepsilon \int_{\Omega} \int_0^t g(t-s) \nabla u(t) \nabla u(s) ds dx \\ &\quad + \varepsilon \int_{\Omega} \int_0^t h(t-s) \nabla v(t) \nabla v(s) ds dx - \varepsilon \int_{\Omega} |u|^{m-2} u_t u dx - \varepsilon \int_{\Omega} |v|^{k-2} v_t v dx \\ &\quad + \varepsilon \int_{\Omega} [uf_1(u, v) + vf_2(u, v)] dx. \end{aligned} \quad (3.10)$$

Consequently, by (3.6) inequality (3.10) can be rewritten as

$$\begin{aligned} L'(t) &\geq -\varepsilon \left(\|\nabla u\|_2^2 + \|\nabla v\|_2^2 \right) + \int_{\Omega} |u|^{m-2} u_t^2 dx + \int_{\Omega} |v|^{k-2} v_t^2 dx \\ &\quad - \varepsilon \int_{\Omega} |u|^{m-2} u_t u dx - \varepsilon \int_{\Omega} |v|^{k-2} v_t v dx + \varepsilon \int_{\Omega} [uf_1(u, v) + vf_2(u, v)] dx \\ &\quad + \varepsilon \int_{\Omega} \int_0^t g(t-s) \nabla u(t) \nabla u(s) ds dx + \varepsilon \int_{\Omega} \int_0^t h(t-s) \nabla v(t) \nabla v(s) ds dx. \end{aligned} \quad (3.11)$$

Therefore, combining with

$$\begin{aligned} \int_{\Omega} \int_0^t g(t-s) \nabla u(t) \nabla u(s) ds dx &= \int_0^t g(s) ds \|u(t)\|_2^2 \\ &+ \int_0^t g(t-s) \int_{\Omega} \nabla u(t) (\nabla u(s) - \nabla u(t)) dx ds. \end{aligned} \quad (3.12)$$

The Cauchy-Schwartz and Young's inequalities allow us to estimate the last term in the right side of (3.12) as follows

$$\int_0^t g(t-s) \int_{\Omega} \nabla u(t) [\nabla u(s) - \nabla u(t)] dx ds \leq \frac{1}{2} (g \circ \nabla u) + \frac{1}{2} \left(\int_0^t g(s) ds \right) \|\nabla u\|_2^2. \quad (3.13)$$

Similarly

$$\begin{aligned} \int_{\Omega} \int_0^t h(t-s) \nabla v(t) \nabla v(s) ds dx &= \int_0^t h(s) ds \|v(t)\|_2^2 \\ &+ \int_0^t h(t-s) \int_{\Omega} \nabla v(t) (\nabla v(s) - \nabla v(t)) dx ds, \end{aligned} \quad (3.14)$$

and

$$\int_0^t h(t-s) \int_{\Omega} \nabla v(t) [\nabla v(s) - \nabla v(t)] dx ds \leq \frac{1}{2} (h \circ \nabla v) + \frac{1}{2} \left(\int_0^t h(s) ds \right) \|\nabla v\|_2^2. \quad (3.15)$$

Inserting (3.12) and (3.14) into (3.11), using (3.13) and (3.15), lead to

$$\begin{aligned} L'(t) &\geq \varepsilon \left(\frac{1}{2} \int_0^t g(s) ds - 1 \right) \|\nabla u\|_2^2 + \varepsilon \left(\frac{1}{2} \int_0^t h(s) ds - 1 \right) \|\nabla v\|_2^2 \\ &+ \varepsilon \int_{\Omega} |u|^{m-2} u_t^2 dx + \varepsilon \int_{\Omega} |v|^{k-2} v_t^2 dx \\ &- \varepsilon \int_{\Omega} [u f_1(u, v) + v f_2(u, v)] dx - \frac{1}{2} \varepsilon (g \circ \nabla u) - \frac{1}{2} \varepsilon (h \circ \nabla v) \\ &- \varepsilon \int_{\Omega} |u|^{m-2} u_t u dx - \varepsilon \int_{\Omega} |v|^{k-2} v_t v dx. \end{aligned} \quad (3.16)$$

To estimate the two last terms in the right hand-side of (3.16), by using the following Young inequality

$$XY \leq \delta X^2 + \delta^{-1} Y^2, \quad X, Y \geq 0, \delta > 0.$$

We have

$$\int_{\Omega} |u|^{m-2} u_t u dx \leq \delta \int_{\Omega} |u|^{m-2} u_t^2 dx + \delta^{-1} \|u\|_m^m. \quad (3.17)$$

Similarly

$$\int_{\Omega} |v|^{k-2} v_t v dx \leq \delta \int_{\Omega} |v|^{k-2} v_t^2 dx + \delta^{-1} \|v\|_k^k. \quad (3.18)$$

Inserting the last two estimates into (3.16), we obtain

$$\begin{aligned}
 L'(t) &\geq \varepsilon \left(\frac{1}{2} \int_0^t g(s) ds - 1 \right) \|\nabla u\|_2^2 + \varepsilon \left(\frac{1}{2} \int_0^t h(s) ds - 1 \right) \|\nabla v\|_2^2 \\
 &\quad + (1 - \varepsilon\delta) \int_{\Omega} |u|^{m-2} u_t^2 dx + (1 - \varepsilon\delta) \int_{\Omega} |v|^{k-2} v_t^2 dx \\
 &\quad - \varepsilon\delta^{-1} \|u\|_m^m - \varepsilon\delta^{-1} \|v\|_k^k + \varepsilon \int_{\Omega} [uf_1(u, v) + vf_2(u, v)] dx \\
 &\quad - \frac{1}{2} \varepsilon (g \circ \nabla u) - \frac{1}{2} \varepsilon (h \circ \nabla v).
 \end{aligned} \tag{3.19}$$

By using (3.1), we get

$$\begin{aligned}
 L'(t) &\geq \varepsilon \left(\frac{1}{2} \int_0^t g(s) ds - 1 \right) \|\nabla u\|_2^2 + \varepsilon \left(\frac{1}{2} \int_0^t h(s) ds - 1 \right) \|\nabla v\|_2^2 \\
 &\quad - \varepsilon\delta^{-1} \|u\|_m^m - \varepsilon\delta^{-1} \|v\|_k^k + (1 - \varepsilon\delta) \int_{\Omega} |u|^{m-2} u_t^2 dx + (1 - \varepsilon\delta) \int_{\Omega} |v|^{k-2} v_t^2 dx \\
 &\quad + \left[2\varepsilon(r+2)H(t) + \frac{2(r+1)}{2} \left(1 - \int_0^t g(s) ds \right) \|\nabla u\|_2^2 \right. \\
 &\quad \left. + \frac{2(r+1)}{2} \left(1 - \int_0^t h(s) ds \right) \|\nabla v\|_2^2 \right. \\
 &\quad \left. + \frac{2(r+1)}{2} (g \circ \nabla u) + \frac{2(r+1)}{2} (h \circ \nabla v) \right] - \frac{\varepsilon}{2} (g \circ \nabla u) - \frac{\varepsilon}{2} (h \circ \nabla v).
 \end{aligned} \tag{3.20}$$

So that

$$\begin{aligned}
 L'(t) &\geq (1 - \varepsilon\delta) \int_{\Omega} |u|^{m-2} u_t^2 dx + (1 - \varepsilon\delta) \int_{\Omega} |v|^{k-2} v_t^2 dx \\
 &\quad - \varepsilon\delta^{-1} \|u\|_m^m - \varepsilon\delta^{-1} \|v\|_k^k + \varepsilon \left(\frac{2(r+2)}{2} - \frac{1}{2} \right) (g \circ \nabla u) \\
 &\quad + \varepsilon \left(\frac{2(r+2)}{2} - \frac{1}{2} \right) (h \circ \nabla v) + 2(r+2)H(t) \\
 &\quad + \varepsilon \left(\frac{2(r+2) - 2}{2} - \frac{1 - 2(r+2)}{2} \int_0^{\infty} g(s) ds \right) \|\nabla u\|_2^2 \\
 &\quad + \varepsilon \left(\frac{2(r+2) - 2}{2} - \frac{1 - 2(r+2)}{2} \int_0^{\infty} h(s) ds \right) \|\nabla v\|_2^2.
 \end{aligned} \tag{3.21}$$

By (3.4), the embedding $L^m(\Omega) \hookrightarrow L^{2(r+2)}(\Omega)$, $L^k(\Omega) \hookrightarrow L^{2(r+2)}(\Omega)$, and using the algebraic inequality

$$z^\nu \leq (z+1) \leq \left(1 + \frac{1}{a} \right) (z+a), \quad \forall z > 0, 0 < \nu \leq 1, a \geq 0, \tag{3.22}$$

we have

$$\begin{aligned}
 \|u\|_m^m &\leq c \|u\|_{2(r+2)}^m \leq c \left(\|u\|_{2(r+2)}^{2(r+2)} \right)^{\frac{m}{2(r+2)}} \\
 &\leq c \left(1 + \frac{1}{H(0)} \right) \left(\|u\|_{2(r+2)}^{2(r+2)} + H(0) \right) \\
 &\leq c_1 \left(\|u\|_{2(r+2)}^{2(r+2)} + H(t) \right).
 \end{aligned}$$

By the Lemma 1 and (3.7), we obtain

$$\|u\|_m^m \leq c_2 \int_{\Omega} [uf_1(u, v) + vf_2(u, v)] dx, \tag{3.23}$$

Similarly

$$\|v\|_k^k \leq c_3 \int_{\Omega} [uf_1(u, v) + vf_2(u, v)] dx. \quad (3.24)$$

Therefore, by (3.23) and (3.24), the inequality (3.21) becomes

$$\begin{aligned} L'(t) &\geq (1 - \varepsilon\delta) \int_{\Omega} |u|^{m-2} u_t^2 dx + (1 - \varepsilon\delta) \int_{\Omega} |v|^{k-2} v_t^2 dx \\ &\quad + 2\varepsilon(r+2)H(t) + \varepsilon a_1 \left(\|\nabla u\|_2^2 + \|\nabla v\|_2^2 \right) \\ &\quad + \varepsilon a_2 [(g \circ \nabla u) + (h \circ \nabla v)] - \varepsilon c_4 \delta^{-1} \int_{\Omega} [uf_1(u, v) + vf_2(u, v)] dx, \end{aligned} \quad (3.25)$$

where

$$\begin{aligned} a_1 &= \min \left(\frac{2(r+2)-2}{2} + \frac{1-2(r+2)}{2} \int_0^{+\infty} g(s) ds, \frac{2(r+2)-2}{2} + \frac{1-2(r+2)}{2} \int_0^{+\infty} h(s) ds \right), \\ a_2 &= \frac{2(r+2)}{2} - \frac{1}{2} \text{ and } c_4 = c_2 + c_3. \end{aligned}$$

Taking

$$0 < a_3 < \min(a_1, a_2),$$

and

$$\begin{aligned} 2H(t) &\geq - \left(\|\nabla u\|_2^2 + \|\nabla v\|_2^2 \right) - [(g \circ \nabla u) + (h \circ \nabla v)] \\ &\quad + \frac{2}{2(r+2)} \int_{\Omega} [uf_1(u, v) + vf_2(u, v)] dx. \end{aligned}$$

We get

$$\begin{aligned} L'(t) &\geq (1 - \varepsilon\delta) \left(\int_{\Omega} |u|^{m-2} u_t^2 dx + \int_{\Omega} |v|^{k-2} v_t^2 dx \right) + \varepsilon(2(r+2) - 2a_3)H(t) \\ &\quad + \varepsilon(a_1 - a_3) \left(\|\nabla u\|_2^2 + \|\nabla v\|_2^2 \right) + \varepsilon(a_2 - a_3) [(g \circ \nabla u) + (h \circ \nabla v)] \\ &\quad + \varepsilon \left(\frac{2a_3}{2(r+2)} - c_4 \delta^{-1} \right) \int_{\Omega} [uf_1(u, v) + vf_2(u, v)] dx \\ &\quad + \varepsilon \left(a_3 \|\nabla u\|_2^2 + \|\nabla v\|_2^2 + (g \circ \nabla u) + (h \circ \nabla v) - \frac{2}{2(r+2)} \int_{\Omega} [uf_1(u, v) + vf_2(u, v)] dx \right). \end{aligned} \quad (3.26)$$

At this point, and for large value of δ and ε small enough such that

$$\frac{2a_3}{2(r+2)} - c_4 \delta^{-1} > 0 \text{ and } 1 - \varepsilon\delta > 0.$$

Noting that

$$2(r+2) - 2a_3 > 0.$$

Then

$$L'(t) \geq M_1 \left(\|\nabla u\|_2^2 + \|\nabla v\|_2^2 + (g \circ \nabla u) + (h \circ \nabla v) + H(t) \right). \quad (3.27)$$

On the other hand, we have

$$L(t) = H(t) + \frac{\varepsilon}{2} \int_{\Omega} u^2 dx + \frac{\varepsilon}{2} \int_{\Omega} v^2 dx + \frac{\varepsilon}{2} \left(\|\nabla u\|_2^2 + \|\nabla v\|_2^2 \right).$$

By Poincaré's inequality, we obtain

$$\int_{\Omega} u^2 dx + \int_{\Omega} v^2 dx \leq c_5 \left(\|\nabla u\|_2^2 + \|\nabla v\|_2^2 \right),$$

so that

$$\begin{aligned} L(t) &\leq M_2 \left(H(t) + \|\nabla u\|_2^2 + \|\nabla v\|_2^2 \right) \\ &\leq M_2 \left(H(t) + \|\nabla u\|_2^2 + \|\nabla v\|_2^2 + (g \circ \nabla u) + (h \circ \nabla v) \right). \end{aligned} \quad (3.28)$$

Then by (3.27) and (3.28), we get

$$L'(t) \geq \xi L(t), \quad (3.29)$$

and

$$L(0) = H(0) + \frac{\varepsilon}{2} \left(\|u_0\|_2^2 + \|v_0\|_2^2 + \frac{\varepsilon}{2} \|\nabla u_0\|_2^2 + \|\nabla v_0\|_2^2 \right) > 0.$$

Finally, a simple integration of (3.29) gives the desired result. \square

Declarations

Conflict of Interest

The authors have no conflicts of interest to declare.

Acknowledgements

The authors wish to thank the anonymous referees for their valuable remarks and careful reading of the proofs presented in this paper.

References

- [1] T. B. BENJAMIN, J. L. BONA, AND J. J. MAHONY, *Model equations for long waves in nonlinear dispersive systems*, Phil. Trans. R. Soc. Lond. A. **272**(1220) (1972), 47–78. DOI
- [2] J. DANG, Q. HU AND H. ZHANG, *Asymptotic stability and blow-up of solutions for the generalized Boussinesq equation with nonlinear boundary condition*. Open J. Math. Anal. **2**(2) (2018), 93–113. DOI
- [3] K. DENG AND H. A. LEVINE, *The role of critical exponents in blow-up theorems: the sequel*, J. Math. Anal. Appl. **243**(1) (2000), 85–126. DOI
- [4] J. ESCERH AND Z. YIN, *Stable equilibria to parabolic systems in unbounded domains*. J. Nonlinear Math. Phys. **11**(2) (2004), 243–255. DOI
- [5] M. ESCOBEDO AND MA. HERRERO, *A semilinear reaction diffusion system in a bounded domain*, Ann. Mat. Pura Appl. **165**(4) (1993), 315–336. DOI
- [6] M. ESCOBEDO AND HA. LEVINE, *Critical blowup and global existence numbers for a weakly coupled system of reaction-diffusion equations*, Arch. Ration. Mech. Anal. **129** (1995), 47–100. DOI

- [7] V. K. KALANTAROV AND O. A. LADYZHENSKAYA, *The occurrence of collapse for quasilinear equations of parabolic and hyperbolic types*, J. Sov. math. **10**(1) (1978), 53–70. [DOI](#)
- [8] G. KARCH, *Asymptotic behavior of solutions to some pseudoparabolic equations*, Mathematical Methods in the Applied Sciences. **20**(3) (1997), 271–289. [DOI](#)
- [9] M. O. KORPUSOV AND A. G. SVESHNIKOV, *Sufficient conditions, which are close to necessary conditions, for the blow-up of the solution of a strongly nonlinear generalized Boussinesq equation*, Comput. Math. Math. Phys. **48**(9) (2008), 1629–1637. [DOI](#)
- [10] H. A. LEVINE, *Some nonexistence and instability theorems for solutions of formally parabolic equations of the form $Pu_t = -Au + \mathcal{F}(u)$* , Arch. Ration. Mech. Anal. **51** (1973), 371–386. [DOI](#)
- [11] S. MESSAOUDI, *Blow-up of solutions of a semilinear heat equation with a memory term*, Abstr Appl Anal. **2** (2005), 87–94. [DOI](#)
- [12] S. MESSAOUDI AND B. SAID-HOUARI, *Global nonexistence of positive initial-energy solutions of a system of nonlinear viscoelastic wave equations with damping and source terms*, J. Math. Anal. Appl. **365**(1) (2010), 277–287. [DOI](#)
- [13] A. OUAOUA, A. KHALDI AND M. MAOUNI, *Stabilisation of solutions for a Kirchhoff type reaction-diffusion equation*, Canad. J. Appl. Math. **2** (2) (2020), 71–80. [URL](#)
- [14] A. OUAOUA AND M. MAOUNI, *Blow-Up, Exponential growth of solution for a nonlinear parabolic equation with $p(x)$ -Laplacian*, Int. J. Anal. Appl. **17**(4) (2019), 620–929. [DOI](#)
- [15] E. PIŞKIN AND F. EKINCI, *Exponential growth of solutions for a parabolic system*, Journal of Engineering and Tecnology. **3**(2) (2019), 29–34.
- [16] N. POLAT, *Blow up of solution for a nonlinear reaction-diffusion equation with multiple nonlinearities*, Int. J. Sci. Technol. **2**(2) (2007), 123–128. [DOI](#)
- [17] A. SVESHNIKOV, G. ALSHIN, A. B. KORPUSOV, AND Y. D. PLETNER, *Linear and nonlinear equations of Sobolev type*, Fizmatlit, Moscow, (2007).
- [18] L. X. TRUONG AND N. VAN Y, *Exponential growth with L_p -norm of solutions for nonlinear heat equations with viscoelastic term*, Appl. Math. Comput. **27**(C) (2016), 656–663. [DOI](#)
- [19] L. X. TRUONG AND N. VAN Y, *On a class of nonlinear heat equations with viscoelastic term*, Comput. Math. Appl. **72**(1) (2016), 216–232. [DOI](#)
- [20] H. ZHANG, J. LU AND Q. Y. HU, *Exponential growth of solution of a strongly nonlinear generalized Boussinesq equation*, Comput. Math. Appl. **68**(12) (2014), 1787–1793. [DOI](#)
- [21] H. W. ZHANG, G. X. ZHANG AND Q. Y. HU, *Global nonexistence of solutions for a class of doubly nonlinear parabolic equations*, Acta Anal. Funct. Appl. **18**(2) (2016), 204–211. [URL](#)