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THE ANALYTICAL SOLUTION OF TELEGRAPH EQUATION OF SPACE-FRACTIONAL ORDER DERIVATIVE BY THE ABOODH TRANSFORM METHOD

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ABSTRACT. In this article, an analytical solution based on the series expansion method is proposed to solve the telegraph equation of space - fractional order (TESFO), namely the Aboodh transformation method (ATM) subjected to the appropriate initial condition. Using ATM, it is possible to find exact solution or a closed approximate solution of a differential equation. Finally, several numerical examples are given to illustrate the accuracy and stability of this method.

1. INTRODUCTION

In the last few decades, fractional calculus found many applications in various fields of physical sciences such as viscoelasticity, diffusion, control, relaxation processes and so on [1]. Suspension flows are traditionally modeled by parabolic partial differential equations. Sometimes they can be better modeled by hyperbolic equations such as the telegraph equation, which have parabolic asymptotic. In particular the experimental data described in [1] seem to be better modeled by the telegraph equation than by the heat equation. The telegraph equation is used in signal analysis for transmission and propagation of electrical signals and also used modeling reaction diffusion. The different type solutions of the fractional telegraph equations have been

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discussed by Momani [2] by using decomposition method, Yildirim [3] by homotopy perturbation method. Our concern in this work is to consider the space-fractional telegraph equations as

$$D_x^{\alpha}u(x, t) = au_t + u_{tt} + bu(x, t) + g(x, t), \ 0 < x < 1$$

where $t \ge 0$, $0 < \alpha \le 2$, a, b are given constants, g(x, t) is given function.

The main objective of this paper is to introduce a new analytical and approximate solution of spatial fractional telegraphic equations using the Aboodh transformation method(ATM), where in [5] authors proposed a Sumudu transformation method (STM) which is used to solve this equation.

2. Preliminary

2.1. Fundamental Properties of Fractional Calculus. In this section we give definitions and some basic results.

Definition 2.1. An Aboodh transform is defined for functions of exponential order. We consider functions in the set F defined by;

$$F = \left\{ f(t) : |f(t)| < Me^{-vt}, \text{ if } t \in [0; \infty[, M, k_1, k_2 > 0; k_1 \le v \le k_2] \right\}$$
(2.1)

For a given function in the set F, M must be finite number and k_1 , k_2 may be finite or infinite with variable v define as $k_1 \le v \le k_2$. Then, the Aboodh transform denoted by the operator A(:) is defined by the integral equation:

$$T(v) = A[f(t)] = \frac{1}{v} \int_{0}^{\infty} f(t)e^{-vt}dt, \quad t \ge 0, \ k_1 \le v \le k_2.$$
(2.2)

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(2.3)

Standard Aboodh transform for some special functions found are given below in Table (2.1).

f(t)	$T\left(v\right) = A\big[f(t)\big]$
1	$\frac{1}{v^2}$
t	$\frac{1}{v^3}$
$t^n, n \ge 1$	$\frac{n!}{v^{n+2}}$
e^{at}	$\frac{1}{v^2 - av}$
sin(at)	$\frac{1}{v(v^2+a^2)}$
cos(at)	$\frac{1}{v^2 + a^2}$
sinh(at)	$\frac{1}{v(v^2 - a^2)}$
$t \cosh(at)$	$\frac{1}{v^2 - a^2}$

TABLE(2.1): Aboodh transform of some functions.

Definition 2.2. The Riemann-Liouville fractional integral of order $\alpha \in R^+$ is defined as

$$D^{-\alpha}f(t) = I^{\alpha}f(t) = \frac{1}{\Gamma(\alpha)} \int_{0}^{t} \frac{f(x)}{(t-x)^{1-\alpha}} dx, \ 0 < \alpha \le 1$$

$$I^{0}f(t) = f(t)$$
(2.4)

Properties of the operator I^{α} can be found in for $\alpha, \beta \ge 0$, and $\gamma \ge -1$, we have:

$$I^{\alpha}I^{\beta}f(t) = I^{\beta}I^{\alpha}f(t) = I^{\alpha+\beta}f(t)$$
$$I^{\alpha}t^{\gamma} = \frac{\Gamma(\gamma+1)}{\Gamma(\alpha+\gamma+1)}t^{\alpha+\gamma}$$

Definition 2.3. The caputo fractional derivative (CFD) operator D_t^{α} of order α is

$$D_t^{\alpha} f(t) = I^{n-\alpha} D^n f(t) = \frac{1}{\Gamma(n-\alpha)} \int_0^t \frac{f^{(n)}(x)}{(t-x)^{1+\alpha-n}} dx, \ x > 0$$
(2.5)

for $n-1 < \alpha \leq n, n \in \mathbb{N}, t > 0$.

Definition 2.4. The Mittage Leffler function $E_{\alpha}(z)$ with $\alpha > 0$, is definite by the following series:

$$E_{\alpha}\left(z\right) = \frac{z^{\alpha}}{\Gamma\left(n\alpha + 1\right)}, \ z \in \mathbb{C}$$

$$(2.6)$$

where $n \in Z^+$, $\alpha \in R^+$.

Definition 2.5. The Aboodh transform $A[D_x^{\alpha}f(x)]$ of the fractional derivative using the Caputo idea of the function is given by:

$$A\left[D_x^{\alpha}f(x)\right] = v^{\alpha}T\left(v\right) - \sum_{k=0}^{n-1} \frac{f^{(k)}(0)}{v^{2-\alpha+k}}$$
(2.7)

It is easy to understand that:

$$A\left[D_t^{\alpha}f(x;t)\right] = v^{\alpha}A\left[f(x;t)\right] - \sum_{k=0}^{n-1} \frac{f^{(k)}(x;0)}{v^{2-\alpha+k}}, \ n-1 < \alpha \le n,$$
(2.8)

Remark 2.1. The Aboodh transform is linear, i.e., if α and β are any constants and f(t) and g(t) are functions defined over the set F above, then

$$A[\alpha f(t) \pm \beta g(t)] = \alpha A[f(t)] \pm \beta A[g(t)].$$

3. PROCEDURE SOLUTION USING ATM FOR SOLVING LINEAR TESFO

We consider the following linear TESFO of the form:

$$D_x^{\alpha} u(x, t) = a u_t + u_{tt} + b u(x, t) + g(x, t), \ 0 < x < 1,$$
(3.1)

$$t \ge 0, \ 0 < \alpha \le 2$$

where g(x, t) is the source term and a, b are constants.

With Initial Condition

$$\frac{\partial^{(r)}u(0,t)}{\partial x^r} = u^{(r)}(0,t)\Big|_{t=0} = f_r(t), r = 0, 1, 2, \dots, n-1.$$
(3.2)

Now applying the AT into Eq(3.1) we have:

$$A[D_x^{\alpha}u(x, t)] = A[au_t + u_{tt} + bu(x, t)] + A[g(x, t)]$$
(3.3)

Substituting Eq(2.8) into Eq(3.3) we get:

$$v^{\alpha}A[u(x;t)]] - \sum_{k=0}^{m-1} \frac{u^{(k)}(0;t)}{v^{2-\alpha+k}} = A[au_t + u_{tt} + bu(x,t)] + A[g(x,t)]$$
(3.4)

$$A[u(x; t)] = \sum_{k=0}^{m-1} \frac{f_k(t)}{v^{2+k}} + v^{-\alpha} A[au_t + u_{tt} + bu(x, t)] + v^{-\alpha} A[g(x, t)]$$
(3.5)

So, according to Aboodh decomposition method (ADM) we can obtain the solution result u(x, t) as

$$u(x, t) = \sum_{n=0}^{\infty} u_n(x, t)$$
(3.6)

Now, substituting Eq(3.6) into Eq(3.5) gives

$$A\left[\sum_{n=0}^{\infty} u_n(x,t)\right] = \sum_{k=0}^{m-1} \frac{f_k(t)}{v^{2+k}} + v^{-\alpha} A\left[a\left(\sum_{n=0}^{\infty} u_n(x,t)\right)_t + \left(\sum_{n=0}^{\infty} u_n(x,t)\right)_{tt} + b\sum_{n=0}^{\infty} u_n(x,t)\right] + v^{-\alpha} A\left[g(x,t)\right]$$
(3.7)

From Eq(3.7) we can define all the coefficients of $u_{n+1}(x, t)$

So we get the zero coefficients $u_0(x, t)$ as:

$$A[u_0(x, t)] = \sum_{k=0}^{m-1} \frac{f_k(t)}{v^{2+k}}$$

The first component $u_1(x, t)$ as:

$$A[u_1(x, t)] = v^{-\alpha} A[a(u_0(x, t))_t + (u_0(x, t))_{tt} + bu_0(x, t) + g(x, t)]$$

Finally the remaining coefficients of $u_{n+1}(x, t)$ can be find in a way like each coefficients is found by using the coming before components.

$$A[u_{n+1}(x, t)] = v^{-\alpha} A[a(u_n(x, t))_t + (u_n(x, t))_{tt} + bu_n(x, t) + g(x, t)], \ n \ge 0.$$

Applying the Aboodh inverse to the above equations yields the following:

$$u_0(x, t) = A^{-1} \left[\sum_{k=0}^{m-1} \frac{f_k(t)}{v^{2+k}} \right]$$

$$u_1(x, t) = A^{-1} \left[v^{-\alpha} A \left[a \left(u_0(x, t) \right)_t + (u_0(x, t))_{tt} + b u_0(x, t) + g(x, t) \right] \right]$$

$$u_{n+1}(x, t) = A^{-1} \left[v^{-\alpha} A \left[a \left(u_n(x, t) \right)_t + \left(u_n(x, t) \right)_{tt} + b u_n(x, t) + g(x, t) \right] \right]$$

So that, the AS $u_n(x, t)$ is given as:

$$u_n(x, t) = \sum_{j=0}^{n-1} u_j(x, t)$$
(3.8)

Such that

$$\lim_{n \to \infty} u_n(x, t) = u(x, t) \tag{3.9}$$

4. Illustrative Examples

In this section we shall test two examples using the ATM to solve the TESFO and the solutions we got it by using the present procedure will be comparing with original ES.

Example 4.1. consider the following homogeneous TESFO

$$D_x^{\alpha}u(x, t) = u_{tt} + u_t + u, \quad x, t \ge 0, \ 0 < \alpha \le 2,$$
(4.1)

$$\begin{cases} u(0, t) = e^{-t}, t \ge 0 \\ u_x(0, t) = e^{-t}, t \ge 0 \end{cases}$$
(4.2)

we appling the AT with (2.8) into (4.1) and (4.2) we get:

$$v^{\alpha}A[u(x;t)] - \sum_{k=0}^{1} \frac{u^{(k)}(0;t)}{v^{2-\alpha+k}} = A[u(x,t)_{tt} + u(x,t)_{t} + u(x,t)]$$
(4.3)

So, we have

$$A[u(x; t)] = \frac{e^{-t}}{v^2} + \frac{e^{-t}}{v^3} + v^{-\alpha}A[u(x, t)_{tt} + u(x, t)_t + u(x, t)]$$
(4.4)

So, according to ADM we can obtain the solution result u(x, t) as

$$u(x, t) = \sum_{n=0}^{\infty} u_n(x, t)$$

substituting (3.6) into (4.4) gives

$$A\left[\sum_{n=0}^{\infty} u_n(x,\,t)\right] = \frac{e^{-t}}{v^2} + \frac{e^{-t}}{v^3} + v^{-\alpha}A\left[\left(\sum_{n=0}^{\infty} u_n(x,\,t)\right)_{tt} + \left(\sum_{n=0}^{\infty} u_n(x,\,t)\right)_{t} + \left(\sum_{n=0}^{\infty} u_n(x,\,t)\right)\right]$$
(4.5)

according to equation (4.5), we can calculate the terms $u_{n+1}(x, t)$

So, we get the coefficients of $u_0(x, t)$ as

$$A[u_0(x; t)] = \frac{e^{-t}}{v^2} + \frac{e^{-t}}{v^3}$$
(4.6)

So, we can use the Aboodh inverse in (4.6), we get

$$u_0(x; t) = A^{-1} \left[\frac{e^{-t}}{v^2} + \frac{e^{-t}}{v^3} \right] = e^{-t} + xe^{-t}$$

and in the same way we calculate the coefficients of $u_1(x,\,t)$

$$A[u_1(x, t)] = v^{-\alpha} A[(u_0(x, t))_{tt} + (u_0(x, t))_t + (u_0(x, t))]$$
(4.7)

Also, we have

$$u_1(x, t) = A^{-1} \left[v^{-\alpha} A \left[(u_0(x, t))_{tt} + (u_0(x, t))_t + (u_0(x, t)) \right] \right]$$
(4.8)

$$\begin{cases} u_{1}(x, t) = A^{-1} \left[v^{-\alpha} A \left[e^{-t} + x e^{-t} \right] \right] \\ = A^{-1} \left[\frac{e^{-t}}{v^{\alpha}} + \frac{e^{-t}}{v^{\alpha+3}} \right] = e^{-t} A^{-1} \left[\frac{1}{v^{\alpha+2}} + \frac{1}{v^{\alpha+3}} \right] \\ = e^{-t} \left(\frac{x^{\alpha}}{\Gamma(\alpha+1)} + \frac{x^{\alpha+1}}{\Gamma(\alpha+2)} \right) \end{cases}$$
(4.9)

We can find the coefficients of $u_n(x, t)$ with the recurrent relation as follows

$$u_{n+1}(x, t) = A^{-1} \left[v^{-\alpha} A \left[(u_n(x, t))_{tt} + (u_n(x, t))_t + (u_n(x, t)) \right] \right], \ \forall n \ge 0$$
(4.10)

Also, we have

$$u_{2}(x, t) = e^{-t} \left(\frac{x^{2\alpha}}{\Gamma(2\alpha+1)} + \frac{x^{2\alpha+1}}{\Gamma(2\alpha+2)} \right)$$

$$\vdots$$

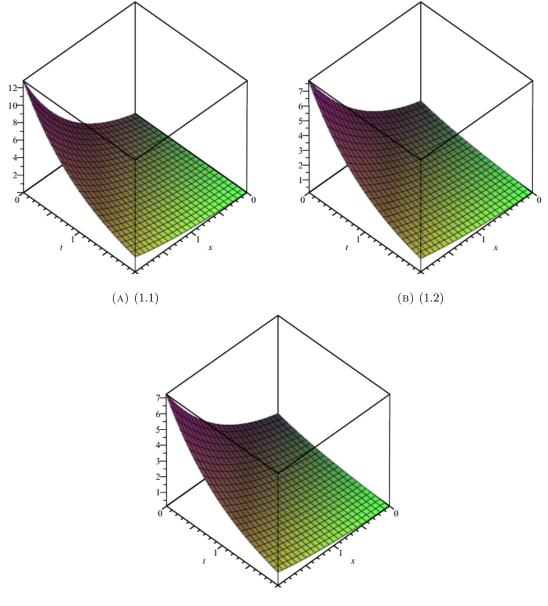
$$u_{n}(x, t) = e^{-t} \left(\frac{x^{n\alpha}}{\Gamma(n\alpha+1)} + \frac{x^{n\alpha+1}}{\Gamma(n\alpha+2)} \right)$$
(4.11)

Finally, we obtain the approximate solution

$$u_n(x, t) = e^{-t} \left(1 + x + \frac{x^{\alpha}}{\Gamma(\alpha+1)} + \frac{x^{\alpha+1}}{\Gamma(\alpha+2)} + \frac{x^{2\alpha}}{\Gamma(2\alpha+1)} + \frac{x^{2\alpha+1}}{\Gamma(2\alpha+2)} + \dots \right)$$
(4.12)

If we put $\alpha = 1$ in (4.12), we can conclude the exact solution

$$u(x, t) = e^{-t} \left(1 + x + x + \frac{x^2}{2!} + \frac{x^2}{2!} + \frac{x^3}{3!} + \frac{x^3}{3!} + \dots \right) = 2e^{-t+x} - e^{-t}$$



(C) (1.3)

FIGURE 1. Comparison between (1.1) the exact solution for $\alpha = 1$ and (1.2), (1.3) the approximative solutions using 4-term of the ATM for $\alpha = 1.7$ and $\alpha = 1.9$ respectively.

Example 4.2. We consider a linear telegraph equation described by

$$D_x^{\alpha}u(x, t) = u_{tt} + 2u_t + u \quad x, t \ge 0, \ 0 < \alpha \le 2,$$
(4.13)

 $with\ initial\ conditions$

$$u(0, t) = e^{-3t}, t \ge 0$$

$$(4.14)$$
 $u_x(0, t) = 2e^{-3t}, t \ge 0$

we appling the AT with (2.8) into (4.13) and (4.14) we get:

$$v^{\alpha}A\left[u(x;\,t)\right] - \sum_{k=0}^{1} \frac{u^{(k)}(0;\,t)}{v^{2-\alpha+k}} = A\left[u(x,\,t)_{tt} + 2u(x,\,t)_t + u(x,\,t)\right]$$
(4.15)

So, we have

$$A[u(x; t)] = \frac{e^{-3t}}{v^2} + 2\frac{e^{-3t}}{v^3} + v^{-\alpha}A[u(x, t)_{tt} + 2u(x, t)_t + u(x, t)]$$
(4.16)

So, according to ADM we can obtain the solution result u(x; t) as

$$u(x, t) = \sum_{n=0}^{\infty} u_n(x, t)$$
(4.17)

substituting (3.6) into (4.16) gives

$$A\left[\sum_{n=0}^{\infty} u_n(x, t)\right] = \frac{e^{-3t}}{v^2} + 2\frac{e^{-3t}}{v^3} + v^{-\alpha}A\left[\left(\sum_{n=0}^{\infty} u_n(x, t)\right)\right)_{tt} + 2\left(\sum_{n=0}^{\infty} u_n(x, t)\right)_{t} + \left(\sum_{n=0}^{\infty} u_n(x, t)\right)\right]$$
(4.18)

according to equation (4.18), we can calculate the terms $u_{n+1}(x, t)$. So, we get the coefficients of $u_0(x, t)$ as

$$A[u_0(x; t)] = \frac{e^{-3t}}{v^2} + 2\frac{e^{-3t}}{v^3}$$
(4.19)

we use the Aboodh inverse in (4.19), we obtain

$$u_0(x; t) = e^{-3t} + 2xe^{-3t}$$

and in the same way we calculate the coefficients of $u_1(x, t)$

$$A[u_1(x, t)] = v^{-\alpha} A[(u_0(x, t))_{tt} + 2(u_0(x, t))_t + (u_0(x, t))]$$
(4.20)

Also, we have

$$u_1(x, t) = A^{-1} \left[v^{-\alpha} A \left[(u_0(x, t))_{tt} + 2 \left(u_0(x, t) \right)_t + (u_0(x, t)) \right] \right]$$
(4.21)

$$\begin{cases} u_{1}(x, t) = A^{-1} \left[v^{-\alpha} A \left[4e^{-3t} + 8xe^{-3t} \right] \right] \\ = e^{-3t} A^{-1} \left[\frac{4}{v^{\alpha+2}} + \frac{8}{v^{\alpha+3}} \right] \\ = 4e^{-3t} \left(\frac{x^{\alpha}}{\Gamma(\alpha+1)} + \frac{2x^{\alpha+1}}{\Gamma(\alpha+2)} \right) \end{cases}$$
(4.22)

We can find the coefficients of $u_n(x, t)$ with the recurrence relation as follows

$$u_{n+1}(x, t) = A^{-1} \left[v^{-\alpha} A \left[(u_n(x, t))_{tt} + 2 (u_n(x, t))_t + (u_n(x, t)) \right] \right]$$
(4.23)

Also, we have

$$\begin{cases} u_{2}(x, t) = A^{-1} \left[v^{-\alpha} A \left[\left(u_{1}(x, t) \right)_{tt} + 2 \left(u_{1}(x, t) \right)_{t} + \left(u_{1}(x, t) \right) \right] \right] \\ = 4A^{-1} \left[v^{-\alpha} A \left[e^{-3t} \left(\frac{4x^{\alpha}}{\Gamma(\alpha+1)} + \frac{8x^{\alpha+1}}{\Gamma(\alpha+2)} \right) \right] \right] \\ = 4e^{-3t} A^{-1} \left[\frac{4}{v^{2\alpha+2}} + \frac{8}{v^{2\alpha+3}} \right] \\ = 16e^{-3t} \left(\frac{x^{2\alpha}}{\Gamma(2\alpha+1)} + \frac{2x^{2\alpha+1}}{\Gamma(2\alpha+2)} \right) \end{cases}$$
(4.24)

we can give the general solution as follow

$$u_n(x, t) = 4^n e^{-3t} \left(\frac{x^{n\alpha}}{\Gamma(n\alpha+1)} + \frac{2x^{n\alpha+1}}{\Gamma(n\alpha+2)} \right)$$
(4.25)

Finally, we obtain the approximate solution

$$u_{n}(x, t) = e^{-3t} \left(1 + 2x + \frac{4x^{\alpha}}{\Gamma(\alpha+1)} + \frac{8x^{\alpha+1}}{\Gamma(\alpha+2)} + \frac{16x^{2\alpha}}{\Gamma(2\alpha+1)} + \frac{32x^{2\alpha+1}}{\Gamma(2\alpha+2)} + \dots \frac{4^{n}x^{n\alpha}}{\Gamma(n\alpha+1)} + 2\frac{4^{n}x^{n\alpha+1}}{\Gamma(n\alpha+2)} + \dots \right)$$
(4.26)

If we put $\alpha = 2$ in (4.26), we can conclude the exact solution

$$u(x, t) = e^{-3t} \left(1 + 2x + \frac{4x^2}{2!} + \frac{8x^3}{3!} + \frac{16x^4}{4!} + \dots \right) = e^{-3t} e^{2x} = e^{-3t+2x}$$
(4.27)

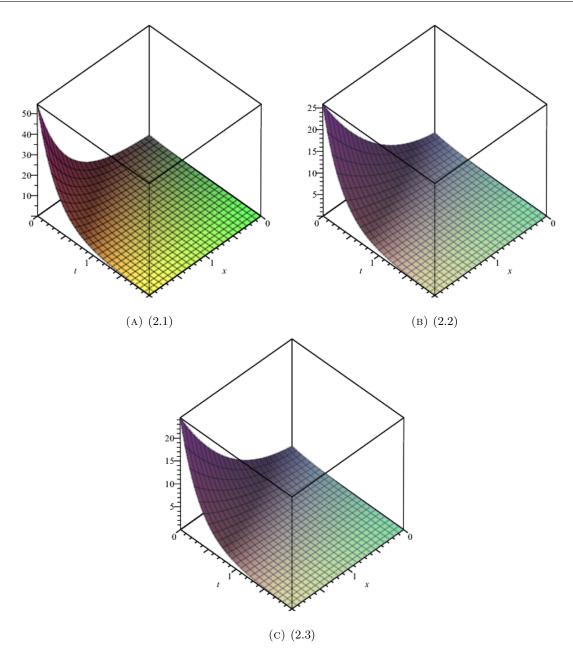


FIGURE 2. Comparison between (2.1) the exact solution for $\alpha = 2$ and (2.2), (2.3) the approximative solutions using 4-term of the ATM for $\alpha = 1.7$ and $\alpha = 1.9$ respectively.

Conclusion. The application of ATM was extended successfully for solving the TESFO. The ATM was clearly very efficient and powerful technique in finding the approximative solution of the proposed equations. In order to check the effectiveness of the introduced procedure, two numerical examples are tested, by comparing the approximative solution with the exact solution.

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